

# Realistic simulations w/ exact chiral symmetry

T. Kaneko for the JLQCD/TWQCD collaborations

<sup>1</sup>High Energy Accelerator Research Organization (KEK)

<sup>2</sup>Graduate University for Advanced Studies

Workshop @ Atami, Feb. 27, 2009

# 1. introduction: JLQCD/TWQCD collaborations

## JLQCD collaboration

- study lattice QCD on super computer system at KEK

KEK: S.Hashimoto, H.Ikeda, TK, M.Matsufuru, J.Noaki, N.Yamada

YITP: H.Ohki, T.Onogi, E.Shintani, T.Umeda

Tsukuba: S.Aoki, K.Takeda, T.Yamazaki

Nagoya: H.Fukaya

Hiroshima: K-I.Ishikawa, M.Okawa

- FY2000: w/ Fujitsu VPP500 (128GFLOPS)

$N_f = 0$  : KS or Wilson/clover:  $B_K, f_{B,D}, B_B, m_s, \dots$

- FY2005: w/ Hitachi SR8000 (1.2TFLOPS)

$N_f = 2$  : clover :  $M_{msn}, M_{brn}, f_{\pi,K}, B_B$

$N_f = 2 + 1$  : clover : **only**  $M_{msn}, M_{brn}, f_{\pi,K}$  (w/ CP-PACS Collab.)

chiral symmetry breaking due to lattice action

⇒ severely limits applications of gauge ensembles!

# 1. introduction: chiral symmetry on the lattice

## chiral symmetry

characterize low-energy dynamics of QCD through SSB

$$M_{\text{msn}}^2 \propto m_q, \dots$$

## price of chiral symmetry breaking

- operator mixing in renormalization :  ~~$B_K$~~
- modified ChPT : ex. Wilson ChPT (Sharpe-Singleton, 1998, ...)  
additional LECs,  $a^2 \ln[m_q]$  singularities  
OK for  $M_{\text{msn}}, f_{PS}$  at NLO; maybe not for others / NNLO

## Nielsen-Ninomiya no-go theorem, 1981

The following properties can NOT hold simultaneously:

- locality
- translational invariance
- no doublers unphys.species, anomaly
- chiral sym. :  $\{\gamma_5, D\} = 0$

$\Leftrightarrow$

## Ginsparg-Wilson fermions, 1982

- GW relation :  
 $\{\gamma_5, D\} = aD\gamma_5 D/m_0$
- modified chiral symmetry  
 $\delta\bar{q} = \bar{q}\gamma_5, \delta q = \gamma_5(1 - aD/2)q$
- correct anomaly, no mixing
- examples of  $D$  .....?

# 1. introduction: dynamical overlap simulations

## overlap fermions (Neuberger, 1998)

$$S_q = \sum \bar{q} D(m) q,$$

$$D(m) = \left(m_0 + \frac{m}{2}\right) + \left(m_0 - \frac{m}{2}\right) \gamma_5 \operatorname{sgn}[H_w(-m_0)]$$

( $m$  = quark mass,  $m_0$  = a tunable parameter)

- satisfies GW relation exactly!
- computationally demanding...

## JLQCD/TWQCD's project

- precise determination of MEs  $\Rightarrow$  test of Standard Model, ...
- current supercomputer system @ KEK (FY2006–FY2010)  
IBM Blue Gene/L + Hitachi SR11K : **59 TFLOPS** in total

$\Rightarrow$  **embark large-scale unquenched simulations with overlap fermion**

in collaboration w/ T.-W.Chiu and T.H.Sheh in Taiwan (TWQCD)

# Outline

## outline

- Introduction
- simulation method
  - implementation; parameters
- fixed topology
- physics results
  - for  $N_f = 2$ :
  - on-going  $N_f = 3$  runs
  - $\langle J_\mu(x) J_\nu(0) \rangle \Rightarrow$  talk by E. Shintani (Kyoto)
- summary

## 2. simulation method

## 2.1 implementation: overlap action

- overlap-Dirac operator  $\ni \text{sgn}[H_W]$

$$D(m) = \left(m_0 + \frac{m}{2}\right) + \left(m_0 - \frac{m}{2}\right) \gamma_5 \text{sgn}[H_W(-m_0)], \quad m_0 = 1.6$$

- $\text{sgn}[H_W] \sim$  polynomial / rational approx.

Zolotarev (min-max) approx:  $\text{sgn}[H_W] = H_W \{p_0 + \sum_{l=1} p_l / (H_W^2 + q_l)\}$

- nested inversion of Dirac operators : very time consuming

$$D(m)^{-1} \ni Dx \ni H_W^{-1}$$

- discontinuity in  $S_q$

$\Rightarrow$  modified HMC (Fodor-Katz-Szabo, 2004) : very time consuming

- locality

$\Leftarrow$  OK, if  $\sigma[H_W]$  has a positive lower bound

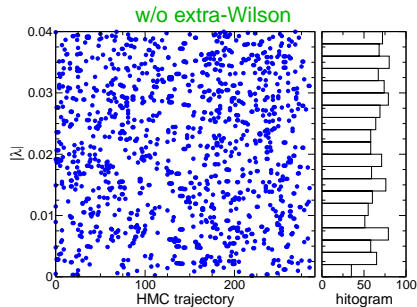
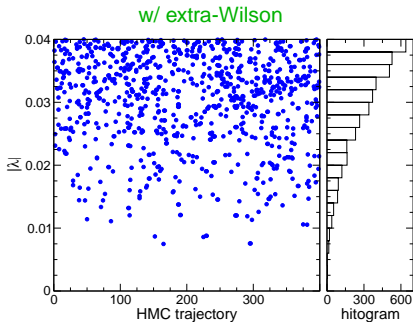
$\Rightarrow$  suppress (near-)zero modes of  $H_W$  by auxiliary Boltzman weight

(JLQCD, 2006)

$$\Delta_W = \frac{\det[H_W(-m_0)^2]}{\det[H_W(-m_0)^2 + \mu^2]}, \quad \mu = 0.2$$

## 2.1 implementation: auxiliary determinant

- $\Delta_W$  : modification of gauge action  $\Rightarrow$  modification of  $a$  dependence
- $\Delta_W$  : suppress (near)zero modes of  $H_W \Rightarrow$  locality, reduced cost



- $\Delta_W$  : fix *global* topology during HMC  $\Rightarrow$  effects should be studied



## 2.1 implementation: other techniques

### other techniques

- multiplications of  $D$  : accuracy of  $\text{sgn}[H_W] = \text{accuracy of chiral sym.}$ 
  - $\sigma[H_W] \Rightarrow [\lambda_{\min}, \lambda_{\text{thrs}}] \cup [\lambda_{\text{thrs}}, \lambda_{\max}]$
  - low modes : Lanczos  $\Rightarrow$  low mode projection
  - high modes : Zolotarev approx. (min-max)
    - $\Rightarrow \langle \nabla_4 A_4 P^\dagger \rangle / \langle P P^\dagger \rangle|_{m=0} \lesssim 0.1(0.1) \text{ MeV}$
- $D$  solver = 5D solver (Edwards et al., 2005) :
  - Schur decomposition ( $D \Rightarrow M_5$ ) + even/odd precondition.  $\Rightarrow$  CPU cost  $\times 1/2$
- HMC algorithm
  - mass preconditioning + multiple time scale MD  $\Rightarrow$  CPU cost  $\times 1/5$   
*Hasenbusch, 2001; Sexton-Weingarten, 1992*
  - no reflection/refraction steps :  $\Rightarrow$  CPU cost  $\times 1/4$
- assembler code for  $H_W$  multiplications :  $\Rightarrow$  CPU cost  $\times 1/3$

## 2.2 parameters

### $N_f = 2$ runs : completed

- Iwasaki-gauge + overlap +  $\Delta_W$
- $\beta = 2.30 \Rightarrow a \approx 0.118(2)$  fm ( $r_0 = 0.49$  fm)
- $16^3 \times 32$  lattice  $\Rightarrow L \simeq 2$  fm
- 6 values of  $m_{ud} \in [m_{s,\text{phys}}/6, m_{s,\text{phys}}]$  for  $Q = 0$
- 10000 HMC trajectories at each  $m_{ud}$
- $Q = -2, -4$  at  $m_{ud} \sim m_{s,\text{phys}}/2$

### $N_f = 2 + 1$ runs : on-going

- $\beta = 2.30 \Rightarrow a \approx 0.107(1)$  fm ( $r_0 = 0.49$  fm)
- $16^3 \times 48$  lattice  $\Rightarrow L \simeq 1.7$  fm
  - 2  $m_s$ 's =  $m_{s,\text{phys}}, 1.3 \times m_{s,\text{phys}}$ ; 5  $m_{ud}$ 's  $\in [m_{s,\text{phys}}/6, m_s]$ ; 2500 traj.
- $24^3 \times 48$  lattice  $\Rightarrow L \simeq 2.6$  fm,  $M_\pi L \gtrsim 3.7$ 
  - $m_s \sim m_{s,\text{phys}}$ ;  $m_{ud} \sim m_{s,\text{phys}}/6, m_{s,\text{phys}}/4$ ; 2500 traj.
- simulations with  $Q \neq 0$  are on-going

## 2.3 measurements: low-mode averaging

### low-mode averaging (LMA)

- 100 low-modes  $\Leftarrow$  Lanczos

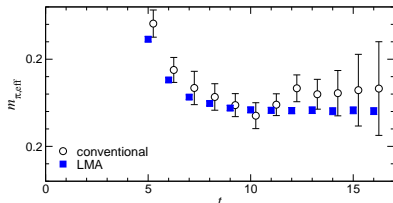
$$D u^{(k)} = \lambda^{(k)} u^{(k)}$$

$$(D^{-1})_{\text{low}} = \sum_k \frac{1}{\lambda_k} u_k u_k^\dagger$$

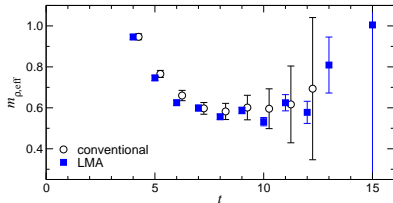
$$C_\Gamma = \left\langle \left\{ (D^{-1})_{\text{low}} + (D^{-1})_{\text{high}} \right\} \Gamma \right. \\ \left. \times \left\{ (D^{-1})_{\text{low}} + (D^{-1})_{\text{high}} \right\} \Gamma \right\rangle \\ = C_{\Gamma,LL} + C_{\Gamma,HL} + C_{\Gamma,LH} + C_{\Gamma,HH}$$

- exact low-mode contrib.  
 $\Rightarrow$  reduce stat. error remarkably
- preconditioning for  $D^{-1}$  solver  
100 modes  $\Rightarrow$   $\times 8$  speed up
- $\rho(\lambda)$ ,  $\chi_t$ , spectrum,  $B_K$ , ...

PS correlator



vector correlator



## 2.3 measurements: all-to-all propagator

### all-to-all quark propagator (TrinLat, 2005)

- high mode  $\Leftarrow$  noise method

$$D x^{(r)} = (1 - P_{\text{low}})\eta^{(r)}$$

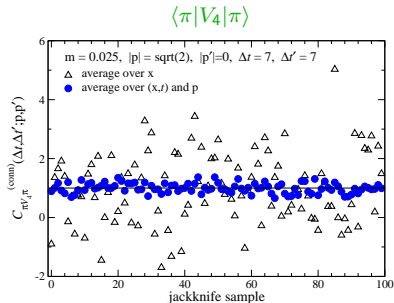
$$(D^{-1})_{\text{high}} = \frac{1}{N_r} \sum_r x^{(r)} \eta^{(r)\dagger}$$

- can construct any correlators w/o small additional CPU costs

- smearing function
- meson momenta

$\Rightarrow$  form factors, hadron decays

- needs large  $N_r$  for  $M_{HL}$ ,  $M_N$ , ...  $\Leftrightarrow O(100)$  TB disk
- pion form factors



### 3. fixed topology

## 3.1 fixed topology

### fixed topology simulations

auxiliary determinant  $\Delta_W \Rightarrow$  fix global topology during HMC : problematic ?

- do NOT suppress local topological fluctuations  
 $\Rightarrow \chi_t$  is calculable from fixed  $Q$  simulations (next slide)
- suppressed by  $1/V$  : small effects ( $\lesssim 1\%$ ) in MEs
- effects can be corrected systematically (Aoki et al., 2007)

$$G_{\theta=0} - G_Q = \frac{G_{\theta=0}^{(2)}}{2\chi_t V} \left( 1 - \frac{Q^2}{\chi_t V} - \frac{c_4}{2\chi_t^2 V} \right) + O(V^{-3})$$

- topological tunneling will be suppressed at  $a \rightarrow 0$   
 $\Rightarrow$  fixing  $Q$  is inevitable for any lattice regularization  
 (w/o modifications of HMC)

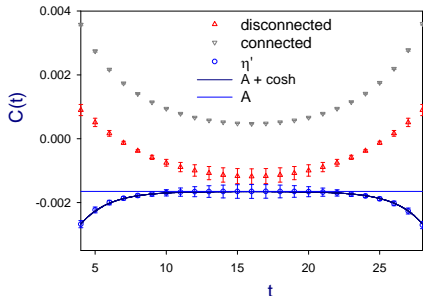
## 3.2 topological susceptibility: determination

### determination from $\eta'$ correlator ( $N_f = 2$ )

$\eta'$  correlator  $\Rightarrow$  a constant term  $\Rightarrow \chi_t$

$$m_q^2 \langle \eta'(t + \Delta t) \eta'(t) \rangle_Q = -\frac{\chi_t}{V} \left( 1 - \frac{Q^2}{\chi_t V} + \frac{c_4}{2\chi_t^2 V} \right) + O(e^{-M_\eta \Delta t}) + O(V^{-3})$$

#### $\eta'$ correlator for $Q=0$



- $Q=0$  : 2nd term vanishes
- 4pt function  $\Rightarrow c_4/(2\chi^2 V)$  : small
- $\eta'$  contamination rapidly damps  
( $M'_\eta > M_\pi$ )  $\Rightarrow$  clear plateau
- 3 sectors  $Q=0, -2, -4$  at  $m_{ud}=0.050$   
 $\Rightarrow$  consistent results for  $\chi_t$

## 3.2 topological susceptibility: chiral behavior

### chiral behavior ( $N_f = 2$ )

LO ChPT (Leutwyler-Smilga, 1992)

$$\chi_t = \frac{\Sigma}{m_u^{-1} + m_d^{-1} + m_s^{-1}} \rightarrow \frac{m_{ud}\Sigma}{2}$$

- expected linear  $m_{ud}$  dependence  
at  $m_{ud} \lesssim m_{s,phys}/2$

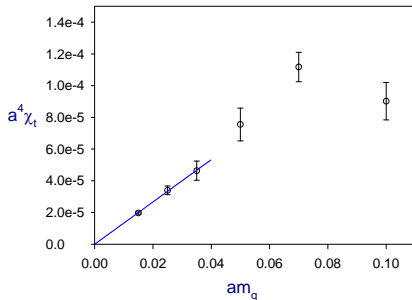
$$\chi_t = c + m_{ud}\Sigma/2 \Rightarrow c = 0.0(1) \times 10^{-5}$$

not clear in previous studies

- ill-determined w/o chiral symmetry  
(index theorem;  $F\tilde{F}$  needs cooling)
- insufficient sampling of topological sectors at small  $m_{ud}$
- provides a determination of  $\Sigma$

$$\Sigma = \{242(5)(10) \text{ GeV}\}^3$$

### $\chi_t$ VS $m_{ud}$



fixed topology strategy :

- introduced to reduce CPU cost...
- an effective way to study topological properties of QCD



## 4. selected results

(mainly from  $N_f = 2$ )

## 4.1 chiral condensate

chiral condensate  $\langle \bar{q}q \rangle \Rightarrow$  SSB of chiral symmetry

### JLQCD/TWQCD's determinations

- topological susceptibility

$$\chi_t = \Sigma m_q / N_f \text{ (degenerate flavors)}$$

$\Leftarrow$  **needs a reliable estimate of  $\chi_t$**

- GMOR relation

$$M_{\text{NG}}^2 = 2B(m_{q,1} + m_{q,2}) + \dots \Rightarrow \Sigma = BF^2$$

$\Leftarrow$  **WChPT modifies the LEC:  $B \rightarrow B(1 - H''/F^2)$**

- eigenvalue distribution in  $\epsilon$ -regime

**comparison w/ ChRMT** (Damgaard-Nishigaki, 2001)

$$\langle \zeta_k \rangle_{\text{ChRMT}} = \langle \lambda_k \rangle_{\epsilon\text{-regime}} \Sigma V$$

- meson correlators in  $\epsilon$ -regime

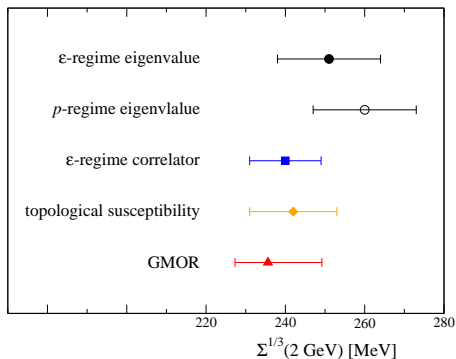
**PS correlator  $\ni \Sigma F$  at NNLO**

$\Leftarrow$  **chiral symmetry + fixed topology in  $\epsilon$ -regime**

#### ChPT in $\epsilon$ -regime

- $M_\pi L \lesssim 1$
- $M_\pi \sim 1/L^2 \sim \epsilon^2$
- modified  $m_q$ -dep.
- $Q$  dependence

## 4.1 chiral condensate

results for  $\langle \bar{q}q \rangle$ 

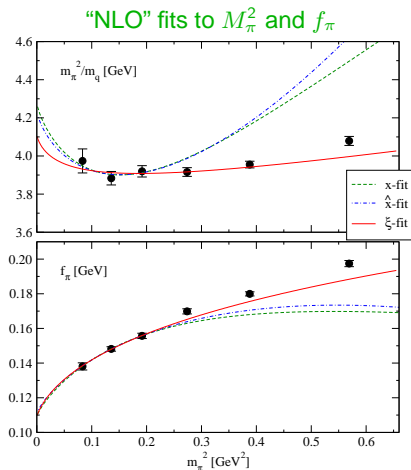
- all results are determined w/  $\lesssim 5\%$  accuracy
- $\langle \bar{q}q \rangle$  from various sources are consistent with each other

⇒ quantitative confirmation of spontaneous chiral symmetry breaking !!

## 4.2 pion mass and decay constant

## pion mass and decay constant

- basic quantities to examine convergence of chiral expansion
  - reliability of LQCD
  - applicability of ChPT
- corrections : few % - few % = small
  - FVC  $\Leftarrow$  NLO ChPT + resummation, (Colangelo et al., 2005)
  - fixed  $Q$   $\Leftarrow$  NLO ChPT w/ fixed  $Q$  (Aoki et al., 2007)
- choice of expansion parameters  
 $x \equiv m_q/F^2$ ,  $\hat{x} \equiv M_\pi^2/F^2$ ,  $\xi \equiv M_\pi^2/F_\pi^2$ 
  - our data at  $m_{ud} < m_{s,phys}/2$  : NLO
  - at  $m_{ud} \lesssim m_{s,phys}$  : NNLO w/  $\xi$   
 $\Leftarrow$  resummation (?)



## 4.2 pion mass and decay constant

low-energy constant

- large deviation in  $l_4$  between NLO and NNLO  
 $\Rightarrow$  NNLO is needed for a reliable estimate of LECs
- results from NNLO fit w/  $\xi$

$$\Sigma = BF^2 = \{235.7(5.0)_{(-2)}^{(+13)} \text{GeV}\}^3$$

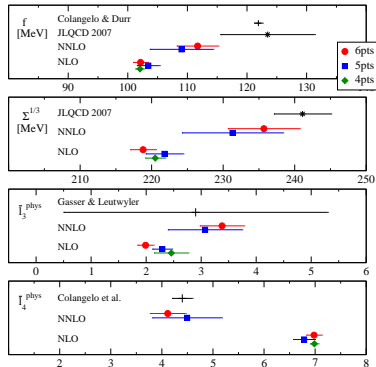
$$F = 79.0(2.5)_{(-0.7)}^{(+4.3)} \text{MeV}$$

$$l_3^r = 3.38(40)_{(-24)}^{(+39)}$$

$$l_4^r = 4.12(35)_{(-30)}^{(+43)}$$

consistent w/ phenomenology (Gasser-Leutwyler, 1984; Colangelo et al., 2001,2004)

## results for LECs



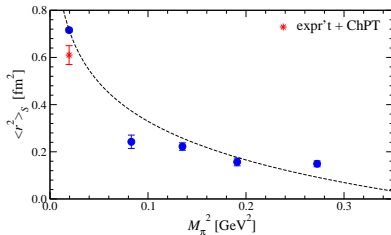
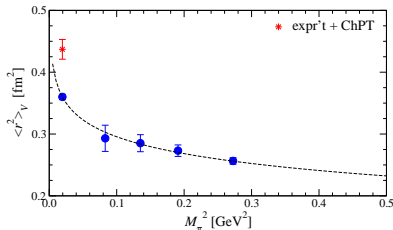
## 4.3 pion form factors

### pion form factors

$$\begin{aligned}\langle \pi(p') | V_\mu | \pi(p) \rangle &= (p + p')_\mu F_V(q^2) \\ \langle \pi(p') | S | \pi(p) \rangle &= F_S(q^2) \\ \langle r^2 \rangle_{V,S} &= 6(\partial F_{V,S}(q^2)/\partial q^2)_{q^2=0}\end{aligned}$$

- next simplest to test convergency
- $\langle r^2 \rangle_S$ 
  - 6 times larger chiral log
  - determination of  $l_4 \Leftrightarrow F_\pi$
  - no experimental process
- measurements w/ **all-to-all**
  - precise estimate of  $F_{V,S}(q^2)$
  - **first calc. of  $F_S$  w/ disc.**
- NLO ChPT fit is not successful
  - $\langle r^2 \rangle_V$  is inconsistent w/ expr't
  - fails to reproduce  $M_\pi^2$  dep. of  $\langle r^2 \rangle_S$

### NLO fits to $\langle r^2 \rangle_{V,S}$



## 4.3 pion form factors

NNLO fit

- 4  $O(p^4)$  LECs, 3  $O(p^6)$  LECs
  - include  $c_V = \partial^2 F_X(q^2)/\partial(q^2)^2$
  - $\bar{l}_3$  : fixed to that from  $M_\pi^2$
  - $\hat{x}$  is used to reduce LECs

- consistent results w/ experiment

$$\langle r^2 \rangle_V = 0.402(22)(36) \text{fm}^2$$

$$\langle r^2 \rangle_S = 0.581(67)(96) \text{fm}^2$$

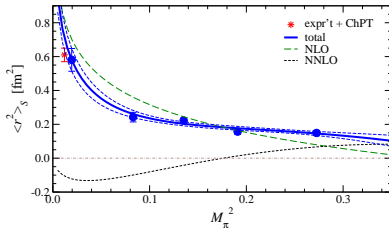
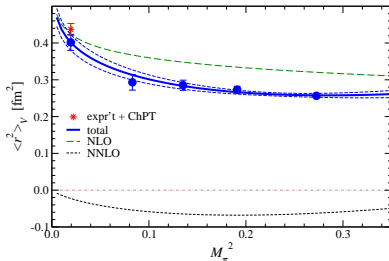
$$\bar{l}_4 = 4.05(55)$$

$$\Leftrightarrow \langle r^2 \rangle_V = 0.437(16) \text{fm}^2 \text{ (expr't)}$$

$$\langle r^2 \rangle_S = 0.61(4) \text{fm}^2 \text{ (expr't)}$$

$$\bar{l}_4 = 4.12(56) \text{ (} F_\pi \text{)}$$

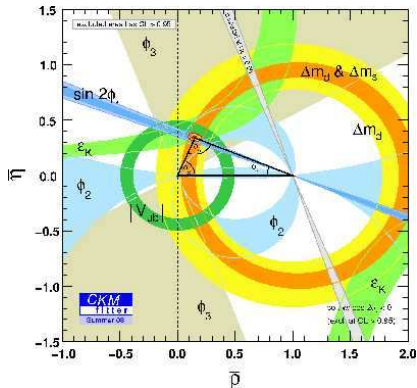
- NNLO contrib. is important in analysis of form factors

NNLO fits to  $F_{V,S}(q^2)$ 

4.4  $B_K$ 

- $B_K \Rightarrow$  tests of Standard Model  
 $B_K \Rightarrow \epsilon_K \Rightarrow$  unitarity triangle
- use of overlap quarks  
 $\Rightarrow$  avoid large error due to op. mixing  
 $VV-AA \Leftrightarrow VV+AA, PP+SS, TT+TT$
- chiral extrap.
  - $M_\pi, F_\pi, \langle r^2 \rangle_{V,S} \Rightarrow$  NNLO ChPT?
  - NNLO formula is not available $\Rightarrow$  use NLO formula w/ NNLO analytic terms

constraints on unitarity triangle





4.4  $B_K$ 

- systematic uncertainties

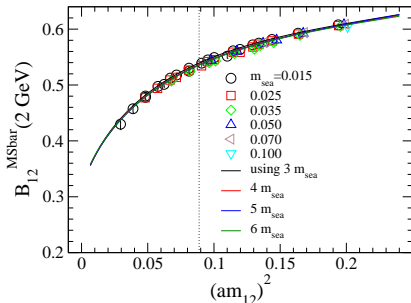
- FVC : 5%  
NLO ChPT + resummation
- fixed  $Q$  : 1.4%  
measurement at  $Q \neq 0$
- scaling violation : 5%  
naive order counting
- $N_f < 3$  : missing  
small ? : cf. RBC :  $N_f = 2 \Leftrightarrow 3$

$$B_K^{\overline{\text{MS}}}(2 \text{ GeV}) = 0.537(4)(40)$$

$$\Leftrightarrow B_K^{\overline{\text{MS}}} = 0.524(10)(28)$$

( $N_f = 3$ ; RBC/UKQCD, 2008)

chiral fit of  $B_K$   
(NLO ChPT + quadratic) fit



expect a better control of sys. error w/  $N_f = 3$  confs at  $M_\pi L \gtrsim 3.7$

## 4.5 baryon quark contents

### $\pi N$ $\sigma$ term

$$\sigma_{\pi N} = m_{ud} \langle N | \bar{u}u + \bar{d}d | N \rangle$$

$$\text{cf. } M_N = M_0 + \sigma_{\pi N} + \dots$$

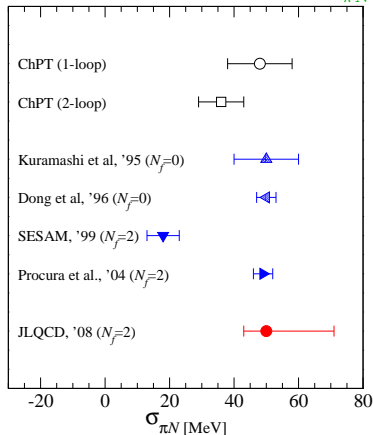
- use Feynman-Hellman theorem

$$\frac{\partial M_N}{\partial m_q} = N_{f,q} m_q \langle N | \bar{q}q | N \rangle$$

- chiral fit of  $M_N$

- $O(p^3)$  BChPT
- $g_A = 1.267$
- ignore some small  $O(M_\pi^{4,5})$  terms
- FVC from BChPT (known up to  $O(p^4)$ )

### ChPT and lattice estimates of $\sigma_{\pi N}$



$$\sigma_{\pi N} = 52(2)_{-7}^{+21} \text{ MeV} \Leftrightarrow \text{consistent w/ previous estimates}$$

## 4.5 baryon quark contents

### y parameter

$$y = \frac{2\langle N|\bar{s}s|N\rangle}{\langle N|\bar{u}u + \bar{d}d|N\rangle}$$

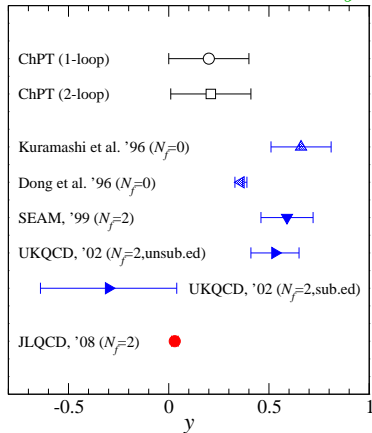
⇒ direct search of dark matter (XMASS, ...)

- chiral fit + Feynman-Hellman
  - $O(p^3)$  PQBChPT
  - $g_A, g_1, \eta_{NN}$  fixed
- chiral symmetry removes a large contamination in “ $\partial M_N / \partial m_{\text{sea}}^{\text{bare}}|_{m_{\text{val}}^{\text{bare}}}$ ” (UKQCD, 2002)

$$\left. \frac{\partial M_N}{\partial m_{\text{sea}}^{\text{phys}}} \right|_{m_{\text{val}}^{\text{phys}}} + \left. \frac{\partial M_N}{\partial m_{\text{val}}^{\text{phys}}} \right|_{m_{\text{sea}}^{\text{phys}}} \frac{\partial m_{\text{val}}^{\text{phys}}}{\partial m_{\text{sea}}^{\text{bare}}|_{m_{\text{sea}}^{\text{phys}}}}$$

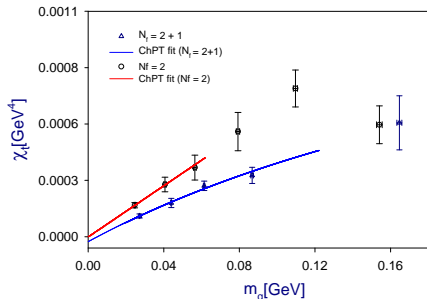
$y = 0.030(16)_{-2}^{+6} \Leftrightarrow$  smaller but natural?; need to extend to  $N_f = 3$

### ChPT and lattice estimates of y



4.6 in  $N_f = 3$  QCD

## topological susceptibility

 $\chi_t$  VS  $m_{ud}$ 

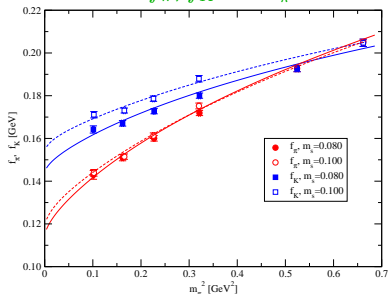
- consistent w/ LO ChPT

$$\chi_t = \Sigma / (2m_{ud}^{-1} + m_s^{-1})$$

- $\Sigma = \{249(4)(2) \text{ MeV}\}^3$

$$\Leftrightarrow \Sigma = \{242(5)(10) \text{ MeV}\}^3 (N_f = 2)$$

## spectroscopy

 $f_\pi, f_K$  VS  $M_\pi^2$ 

- NNLO ChPT fits w/ NPR

$$m_{ud}^{\overline{MS}}(2 \text{ GeV}) = 3.80(7) \text{ MeV}$$

$$m_s^{\overline{MS}}(2 \text{ GeV}) = 115(2) \text{ MeV}$$

$$f_\pi = 121.5(4.1) \text{ MeV}$$

$$f_K = 148.3(4.7) \text{ MeV}$$

$$f_K / f_\pi = 1.220(10)$$

## summary

- JLQCD/TWQCD's dynamical overlap project:
  - advantage : **chiral symmetry**
    - essential to calculate general matrix elements (4 quark op., multi. bilinear, ...)
    - study chiral behavior based on continuum ChPT
  - drawback : **not on physical point**  $\Leftrightarrow$  PACS-CS (*talk by Ukita*)
    - $\Leftrightarrow$  quantities w/ poor convergency of chiral expansion
- future possibilities w/ ensembles at ( $m_{ud} \gtrsim m_{s,phys}/6$ ,  $M_\pi L \gtrsim 3.7$ )
  - current KEK supercomputer system : till FY2010
    - $\Rightarrow$   $Q \neq 0$ ,  $\epsilon$ -regime, degenerate  $N_f = 3$ , high  $T$ , ...
  - what can be done for hyperon physics ?
    - single body problem : may be OK
      - matrix elements ( $\Leftarrow$  overlap), disconnected diagram ( $\Leftarrow$  all-to-all), ...
    - baryon interaction ?
      - $M_\pi L \gtrsim 4$  is safe for multi-baryon system?
      - chiral symmetry plays important roles ?