



Photoproduction of Φ and $\Lambda(1520, 1405)$



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Introduction

Nucleon (N) and hyperon (Y) resonance searches

- Insight to hadron structures and production mechanisms
- Completion of hadron spectrum
- Strangeness physics from Y-resonances
- Exotic baryon searches: pentaquark baryon

Experiments

- LEPS, Tohoku-LNS, CLAS, GRAAL, BONN-ELSA etc..
- η , K and ϕ photo- and electro-productions
- $N^*(1535)$, $\Sigma^*(1385)$, $\Lambda^*(1405)$, $\Lambda^*(1520)$ etc..
- Exotics: $N^*(1675)$ from LNS and GRAAL, and $\Theta(1540)$ from LEPS

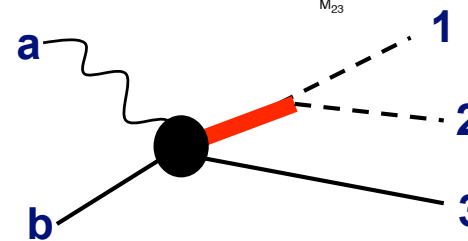
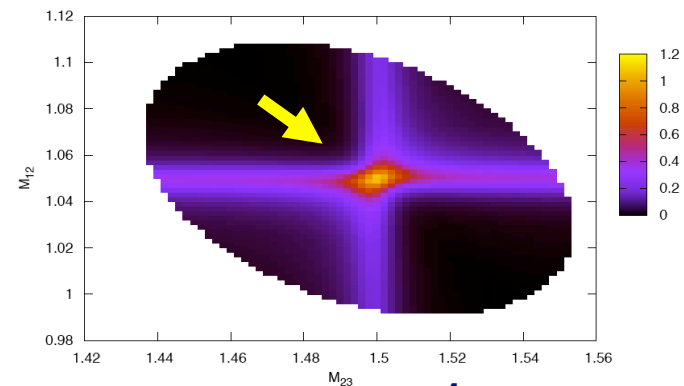
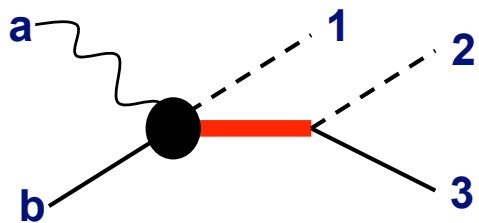
Introduction

Theories

- Coupled channel methods: dynamical generation of resonances
- Effective Lagrangian methods: resonances as element fields
- Baryonic models: Skyrmion with excitation, CQSM, etc..

Resonances via three-body final state

- Three-body final state LIPS
- Dalitz-plot analyses for $a, b \rightarrow 1, 2, 3$
- Resonances appear as strips



Introduction

Interference between resonances

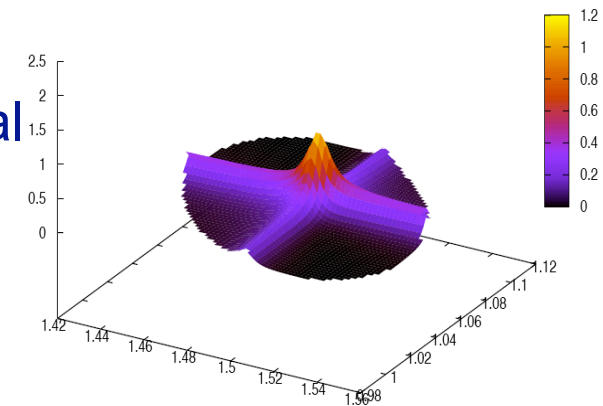
- Constructive interference \rightarrow enhancing signal
- Sensitive to relative phase, strength, etc..

Low statistics in exotic baryon search

- Small coupling strength of $\Theta \rightarrow KN \sim 1 \text{ MeV}$
- $\sigma(\gamma N \rightarrow K\Theta) \sim$ a few nb: hard to be seen in production processes
- Strong constructive interference on Dalitz plot!? [Amarian et al. hep-ph/0612150](#)

False signals

- Resonance-like structure generated from kinematical reasons
- Resonance contribution + phase space



Introduction

A model calculation to test these topics (interference & false signal)

Basic studies for further researches on resonances

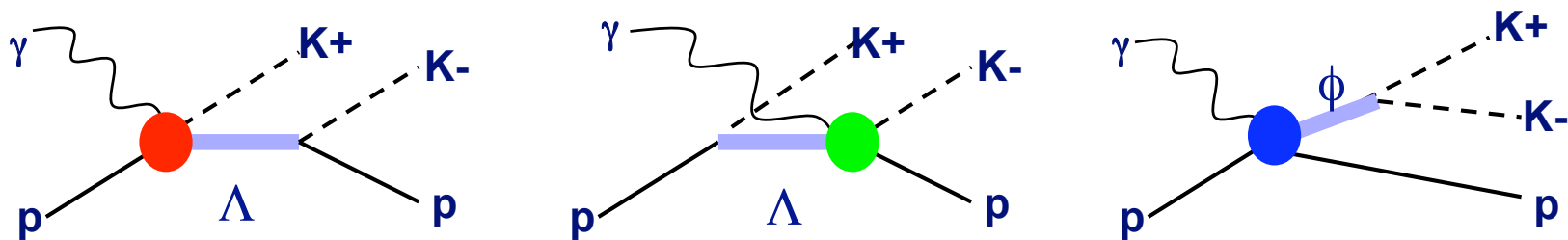
Theoretical method

- Considering $\gamma p \rightarrow K^+ K^- p$ containing $\Lambda(1520)$ and $\phi(1020)$
- Considering the LEPS experimental setup: K^+ only in forward angle
- Employing an effective Lagrangian method
- Form factors in a gauge-invariant manner
- Finally, plotting Dalitz plots

Theoretical framework

Feynman diagrams

- For the interference pattern: $\Lambda(1520)=\Lambda^*$ and $\phi(1020)$ contribution only



- **Red** blob contains s-,t-,u-channels and contact term: gauge invariant
- **Green** blob contains s-,t-,u-like channels, contact term: gauge invariant
- **Blue** blob contains $F_{\mu\nu}$: gauge invariant itself and pomeron effect
- Photon- Λ^* coupled diagram ignored: neutral and no data for κ for Λ^*
- K^* -exchange contribution ignored: $g(K^*N\Lambda^*) \ll 1$

Theoretical framework

Effective Lagrangians

- Λ^* contribution ($\Lambda^*=B^*$)

$$\mathcal{L}_{\gamma KK} = ie_K [(\partial^\mu K^\dagger)K - (\partial^\mu K)K^\dagger] A_\mu + \text{h.c.},$$

$$\mathcal{L}_{KNB^*} = \frac{g_{KNB^*}}{M_{B^*}} \bar{B}^{*\mu} \partial_\mu K \Gamma_5 \gamma_5 N + \text{h.c.},$$

$$\mathcal{L}_{\gamma KNB^*} = -i \frac{e_N g_{KNB^*}}{M_{B^*}} \bar{B}^{*\mu} A_\mu K \Gamma_5 \gamma_5 N + \text{h.c.},$$

- Rarita-Schwing vector-spinor fields for Λ^*
- Λ^* (spin-3/2) propagator \sim spin-1/2 propagator for low-energy region
- $g(KN\Lambda^*)$ obtained from exp. data

Theoretical framework

Effective Lagrangians

- ϕ contribution

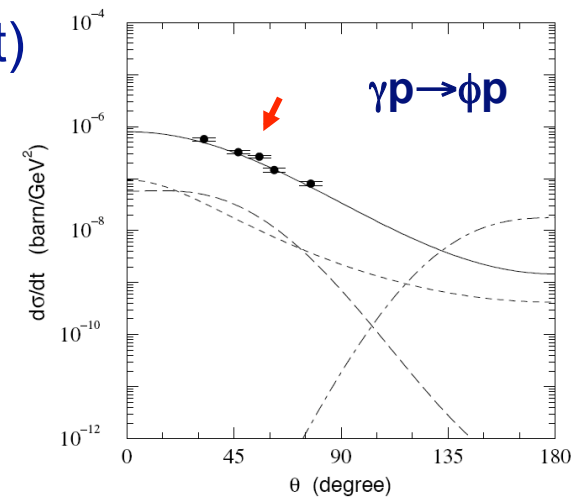
$$\mathcal{L}_{\phi KK} = ig_{KK\phi} \phi^\mu [(\partial_\mu K^\dagger)K - (\partial_\mu K)K^\dagger] + \text{h.c.},$$

$$\mathcal{L}_{\gamma\phi NN} = \frac{\mathcal{F}(s, t)}{M_\phi M_N^2} \bar{N} \partial_\mu \phi_\nu F^{\mu\nu} N + \text{h.c.},$$

- Pomeron contribution taken into account in $F(s, t)$
- Reproduction of strong forward scattering

A.I. Titov et al PRC58,2429

- Model parameters fitted by data



Theoretical framework

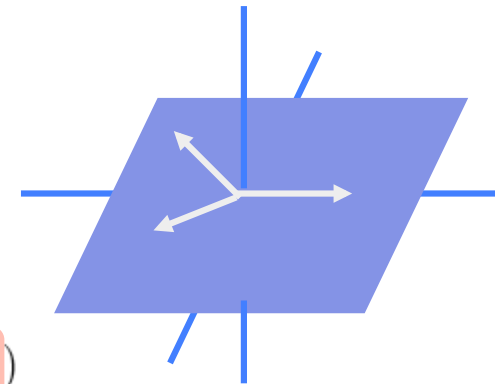
Three-body LIPS and double-differential cross section

- LEPS kinematics simplifies computation significantly: $\mathbf{k}_K/|\mathbf{k}_K| \sim \mathbf{k}_\gamma/|\mathbf{k}_\gamma|$
- Inner-products, $\mathbf{k}_K \cdot \boldsymbol{\varepsilon}_\gamma$ and $\mathbf{k}_p \cdot \boldsymbol{\varepsilon}_\gamma$ become zero in CM frame

$$\frac{d^2\sigma}{dM_{KN_f} d\cos\theta_K} = \int \frac{M_{K\bar{K}} M_{\bar{K}N_f} dM_{K\bar{K}}}{128\pi^3 E_{\text{cm}}^2 E_\gamma} \left[\frac{1}{4} \sum_{\text{spin}} \sum_{\text{pol}} |\mathcal{M}_{\gamma N_i \rightarrow K\bar{K}N_f}|^2 \right]$$

Definition of relevant momenta

$$\begin{aligned} p_\gamma &= (E_\gamma, \mathbf{p}_\gamma) = (|\mathbf{p}_\gamma|, 0, 0, |\mathbf{p}_\gamma|), \\ p_{N_i} &= (E_N, \mathbf{p}_{N_i}) = ([M_N^2 + \mathbf{p}_\gamma^2]^{1/2}, 0, 0, -|\mathbf{p}_\gamma|), \\ p_K &= (E_K, \mathbf{p}_K) = ([M_K^2 + \mathbf{p}_K^2]^{1/2}, 0, 0, |\mathbf{p}_K|), \\ p_{\bar{K}} &= (E_\pi, \mathbf{p}_{\bar{K}}) = p_\gamma + p_{N_i} - p_K - p_{N_f}, \\ p_{N_f} &= (E_\Sigma, \mathbf{p}_{N_f}) = ([M_\Sigma^2 + \mathbf{p}_{N_f}^2]^{1/2}, 0, |\mathbf{p}_{N_f}| \sin\theta_{N_f}, |\mathbf{p}_{N_f}| \cos\theta_{N_f}) \\ E_K &= \frac{s + M_K^2 - M_{\bar{K}N_f}^2}{2\sqrt{s}}, \quad E_{N_f} = \frac{s + M_{N_f}^2 - M_{K\bar{K}}^2}{2\sqrt{s}} \end{aligned}$$



Theoretical framework

Conditions for the Dalitz plot

- Maximum and minimum, confining Dalitz region

$$M_K + M_{\bar{K}} \leq M_{K\bar{K}} \leq \sqrt{s} - M_{N_f},$$

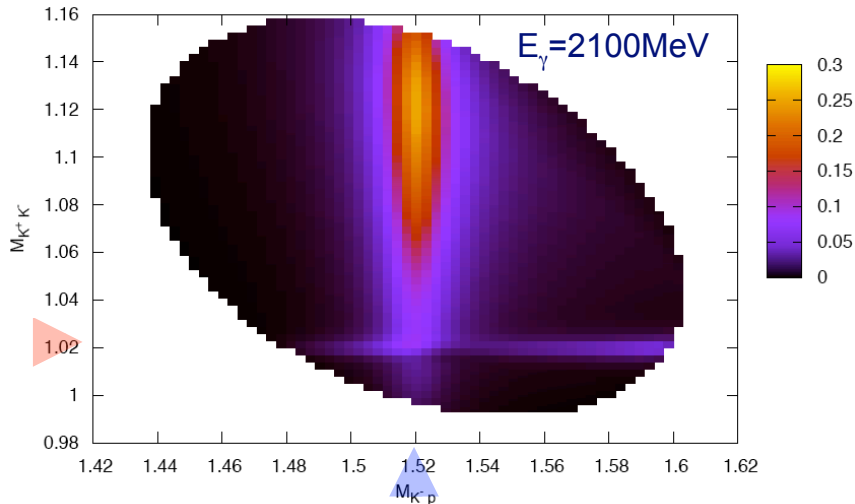
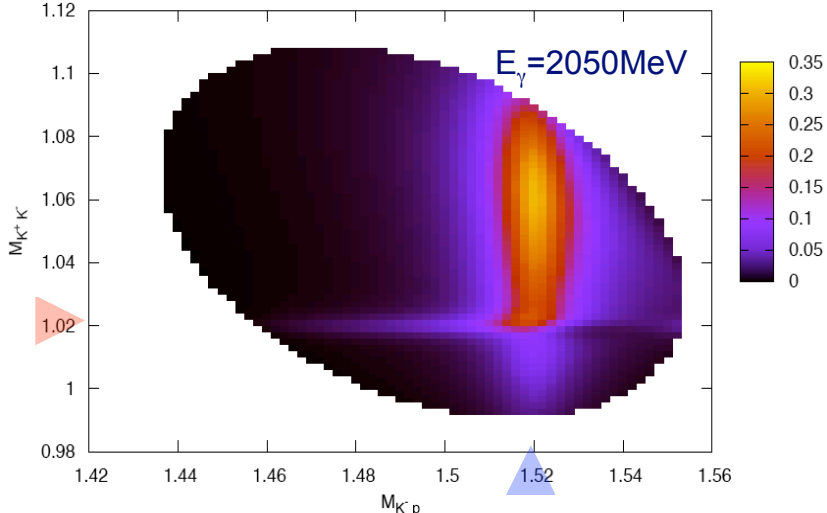
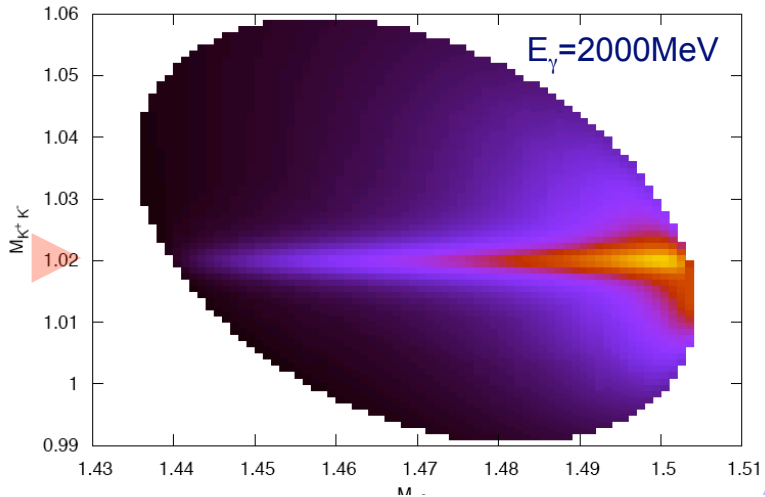
$$M_{\bar{K}} + M_{N_f} \leq M_{\bar{K}N_f} \leq \sqrt{s} - M_K.$$

- Angle between final \mathbf{k}_p and z-axis

$$-1 \leq \cos \theta_{N_f} = \frac{(\sqrt{s} - E_K - E_{N_f})^2 - M_{\bar{K}}^2 - |\mathbf{p}_{\bar{K}}|^2 - |\mathbf{p}_{N_f}|^2}{2|\mathbf{p}_{\bar{K}}||\mathbf{p}_{N_f}|} \leq 1$$

Numerical results: interference

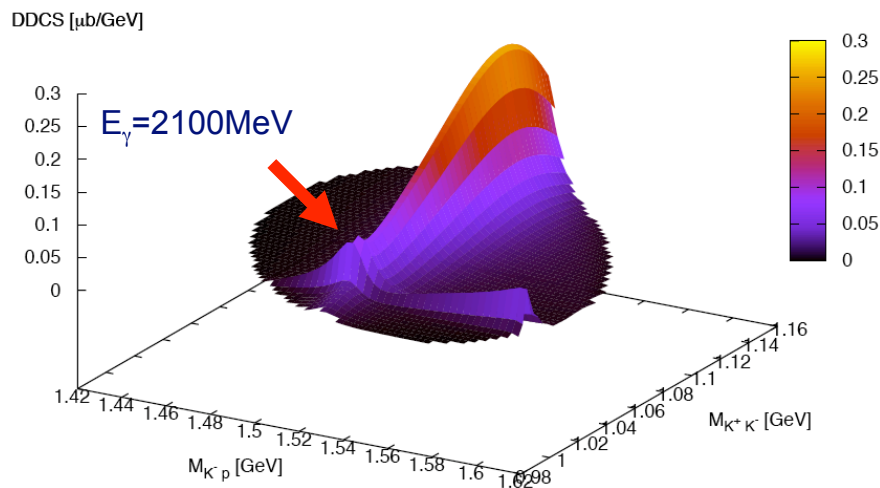
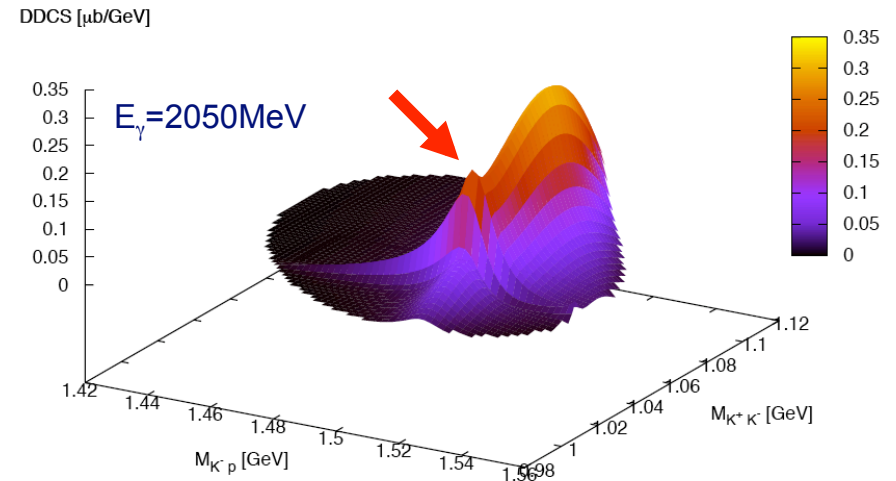
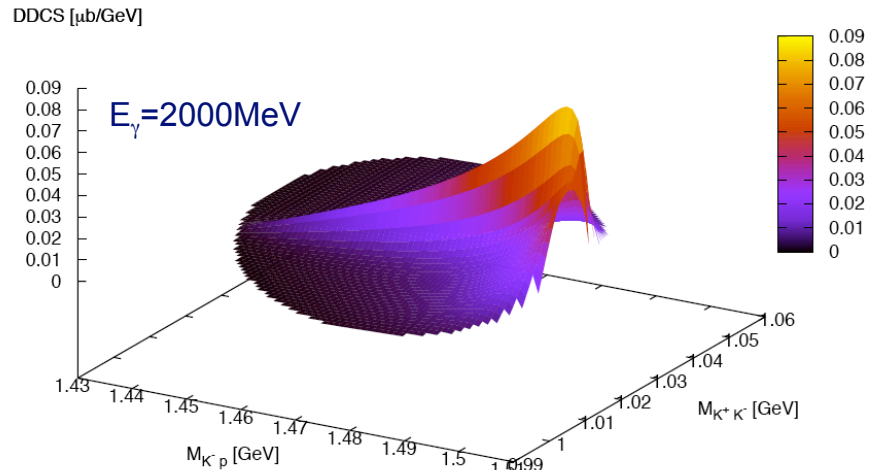
Dalitz plots for $\gamma p \rightarrow K^+ K^- p$



- Triple-differential σ [$\mu\text{b}/\text{GeV}$]
- Small interference effect

Numerical results: interference

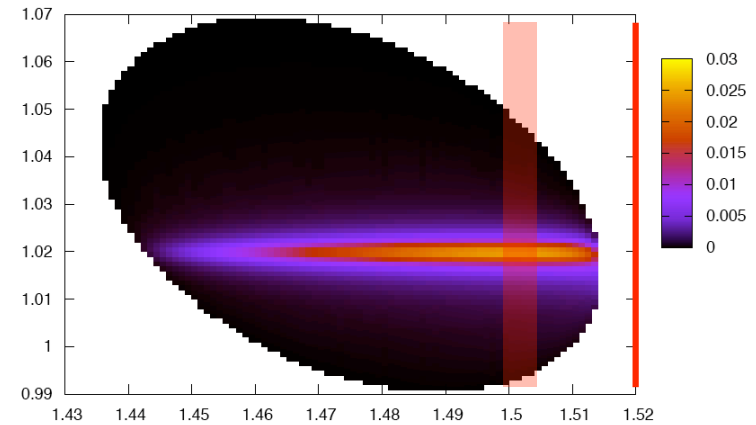
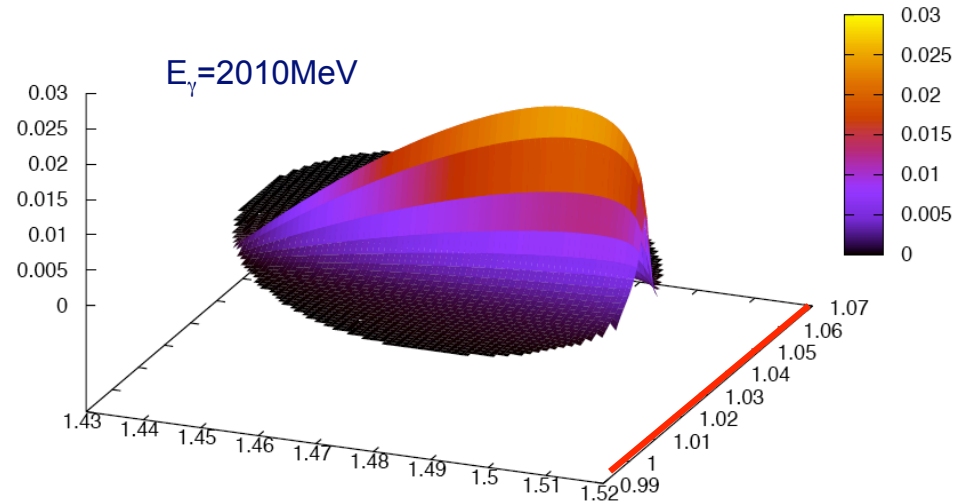
Dalitz plots for $\gamma p \rightarrow K^+ K^- p$



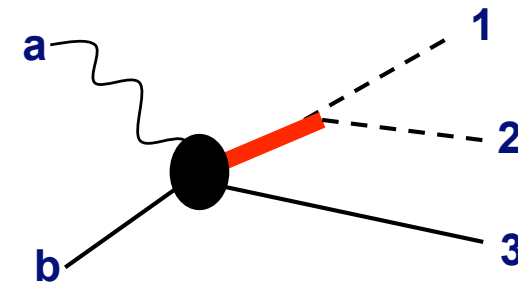
- Small interference effect
- $\sigma_\phi < \sigma_{\Lambda(1520)}$
- Insensitive to relative phase

Numerical results: false signal

Dalitz plots for $\gamma p \rightarrow K^+ K^- p$ with only ϕ contribution



- B.G. contribution + phase space
- False signal at low-energy region
- Depending on experimental resolution?



Summary and conclusion

Studies on three-body final state using Dalitz plot

$\gamma p \rightarrow K^+ K^- p$ containing $\Lambda(1520)$ and $\phi(1020)$

Effective Lagrangian method employed

Interference effect

- Sensitive to strength and/or phase
- Negligible effects on $\gamma p \rightarrow K^+ K^- p$ due to $\sigma_\phi < \sigma_{\Lambda(1520)}$
- If, in $\gamma n \rightarrow K^+ K^- n$, $\sigma_\phi \sim \sigma_\Theta$ at low energy region, the effect becomes visible

False signal

- B.G. contribution + phase space without resonances
- Possible for low-energy region with coarse resolution
- Sensitive to strength and/or phase

Theoretical status

$\Lambda(1405)$ KN bound state rather than a uds color singlet!? cf) $N^*(1535)$

Energetic activities on $\Lambda(1405)$ studies

- Two-pole structure Spain and Japan groups
- Different microscopic structure: larger size Sekihara et al, PLB** [arxiv:0803.4068 [nucl-th]]
- Studies on the origin of s-wave resonance Hyodo et al, PRC78,025203
- Three-body problems KKN-molecular state D.Jido et al, arXiv:0806.3601 [nucl-th]

$\Lambda(1405)$ nature via meson and photon induced reactions

- $\gamma N \rightarrow K\Lambda(1405)$ Williams et al, PRC43,452
- Production mechanism SiN et al, arXiv:0806.4029 [hep-ph]

Focusing on the photoproduction

Experimental status

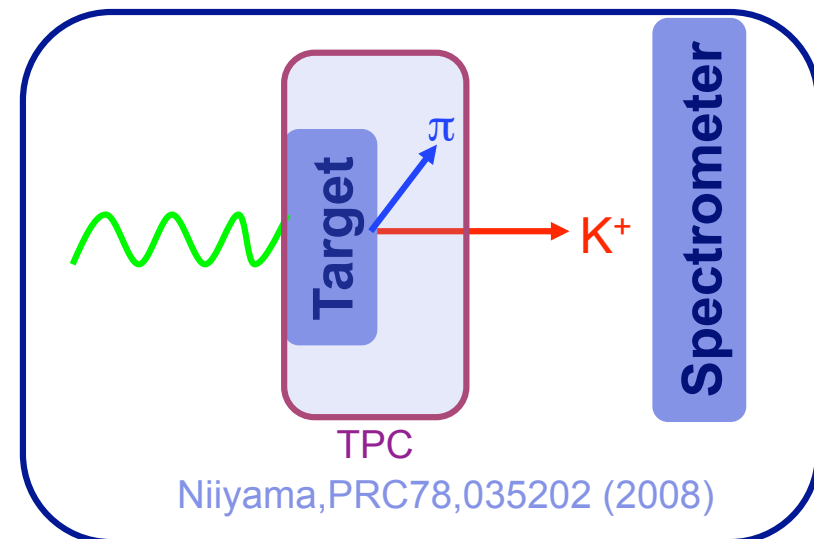
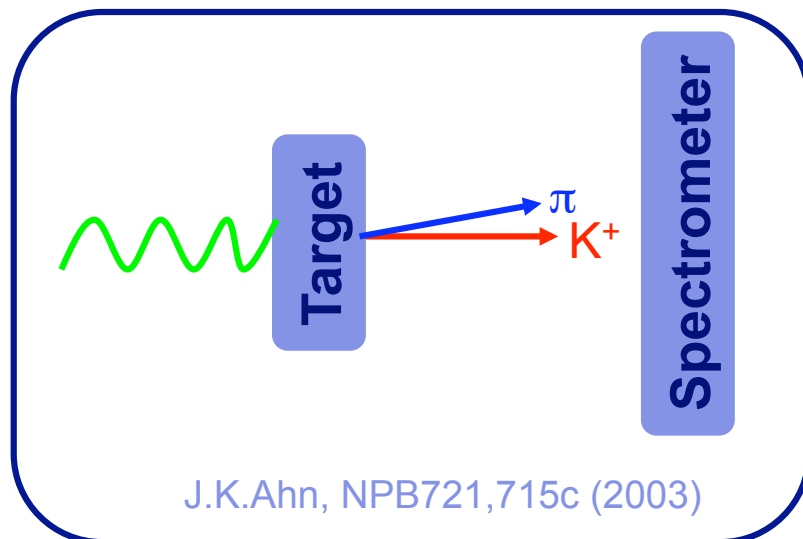
Meson-nucleon scattering

- K^- -N scattering Hepp NPB115,82, Hermingway NPB253,742, Prakhov PRC70,034605
- π^- -p scattering Thomas NPB56,15

Photon-nucleon scattering

- γ -p scattering (by LEPS) J.K.Ahn NPB721,715c, Niiyama et al, PRC78,035202

Main difference between the two photoproduction data

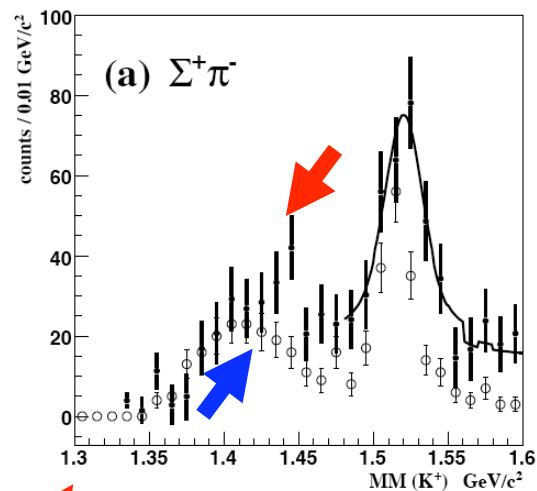


Experimental status: angle dependence

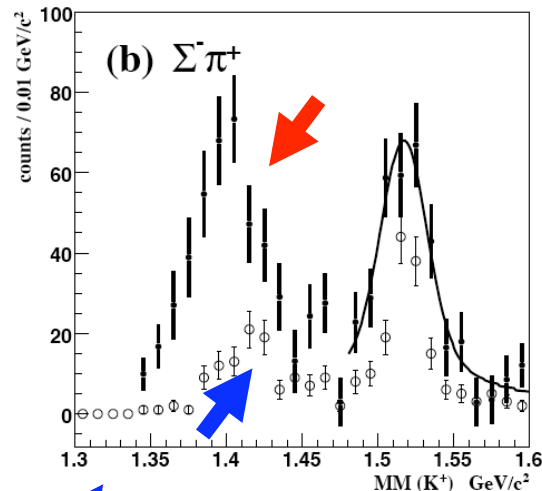
Pion measured in wider angle by TPC

Angle-dependent production becomes critical in the new experiment

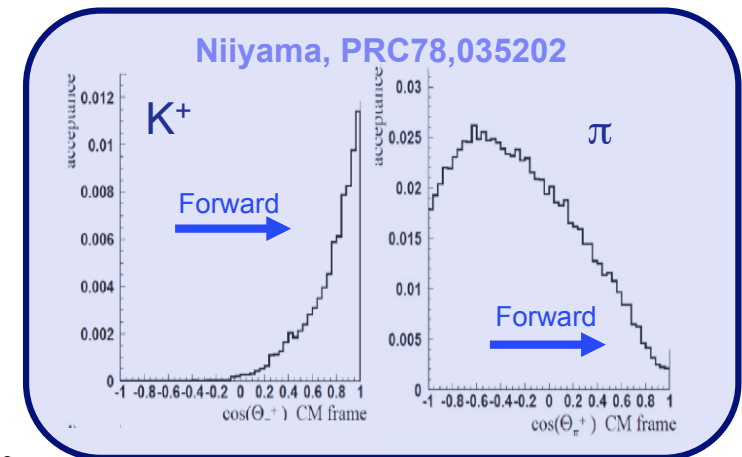
K^+ missing mass in $\gamma p \rightarrow K^+ X$



↗ Niiyama et al, PRC78,035202



↘ J.K.Ahn, NPA721,715c

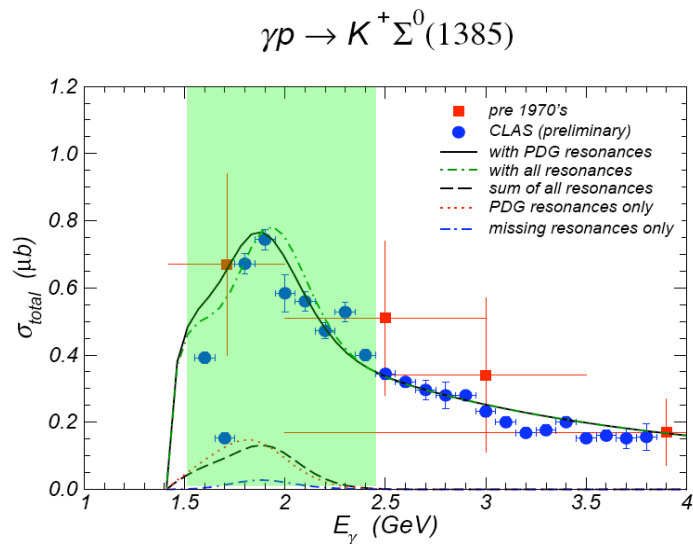
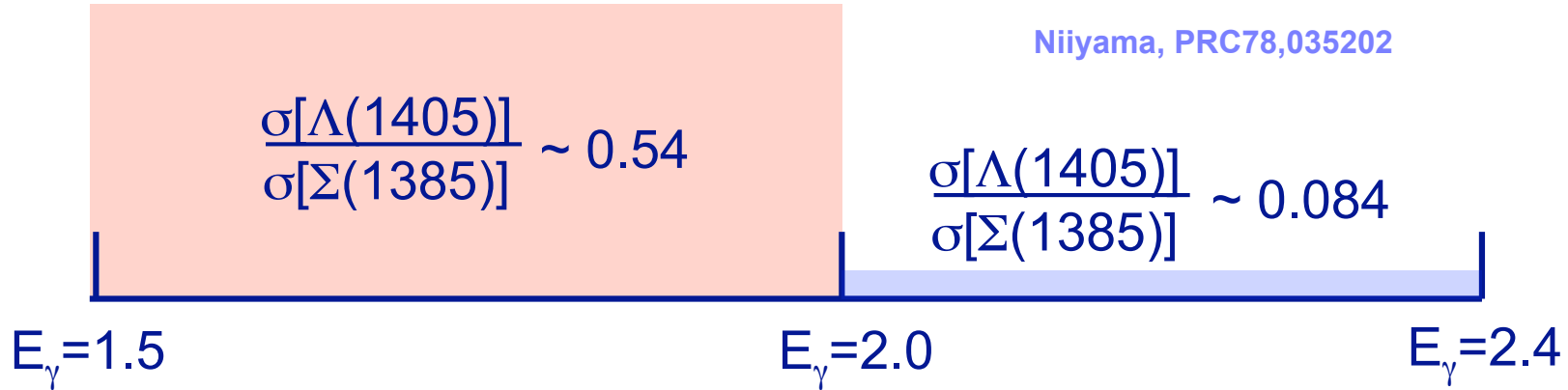


Significant difference between $\pi^+\Sigma^-$ and $\pi^-\Sigma^+$ channels

Indicating sizable p-wave contribution in $\Lambda(1405)$ photoproduction!?

Experimental status: energy dependence

$\gamma p \rightarrow K^+ \Lambda(1405)$ cross section estimated in the new exp.



Y.Oh, PRC77,045204

$\sigma[\Sigma(1385)] \sim 0.6 \mu\text{b}$ for Rg.1
 $\sigma[\Sigma(1385)] \sim 0.5 \mu\text{b}$ for Rg.2

$\sigma[\Lambda(1405)] \sim 0.3 \mu\text{b}$ for Rg.1
 $\sigma[\Lambda(1405)] \sim 0.04 \mu\text{b}$ for Rg.2

Factor ~ 8 difference!

Theoretical analysis

Possible p-wave contribution in $\Lambda(1405)$ mass spectrum

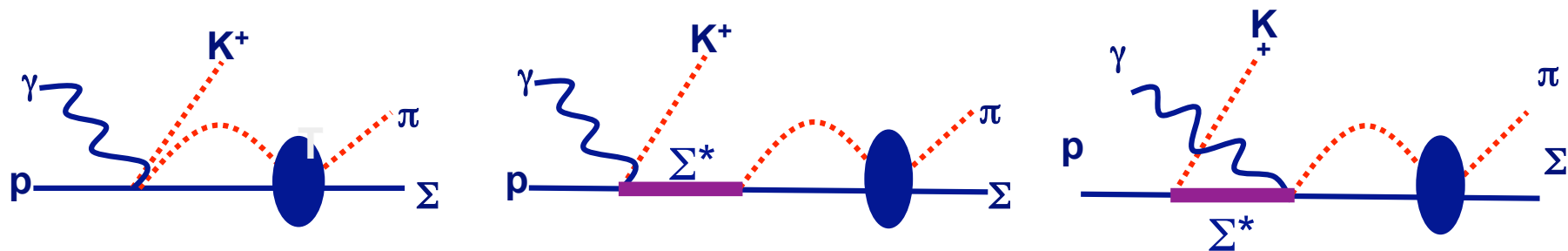
Employing the ChUM with explicit p-wave contribution from $\Sigma^*(1385)$

$\Lambda(1405)$ photoproduction with s-wave by Nacher et al. [Nacher,PLB455,55](#)

Considering the LEPS experimental setup: K^+ in forward angle ($\theta_K \sim 0$)

$$\frac{d^2\sigma}{dM_{\pi\Sigma} d\cos\theta_K} = \int \frac{M_{K\pi} M_{\pi\Sigma} dM_{K\pi}}{128 \pi^3 s^{3/2} E_\gamma} |\overline{\mathcal{M}}|^2$$

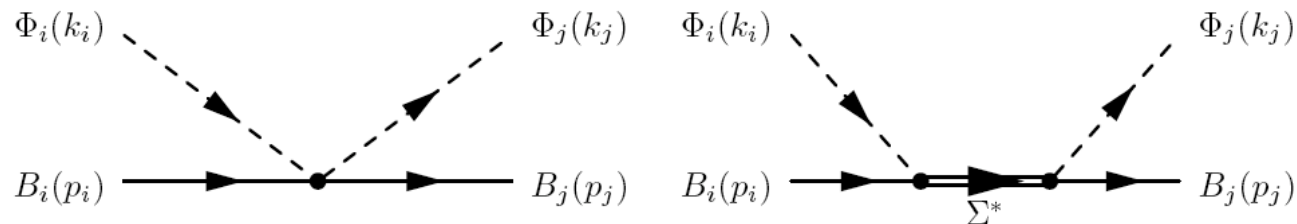
Diagrams considered



[Nacher,PLB455,55](#)

Theoretical analysis

$\Phi B \rightarrow \Phi B$ amplitude (●) including (s,p)-wave contributions Jido,PRC66,055203



Lagrangians for the interaction vertices

$$\mathcal{L}_{\text{WT}} = -\frac{iC_{ij}}{4f^2} \bar{B}_j [\Phi_j(\not{\partial}\Phi_i) - \Phi_i(\not{\partial}\Phi_j)] B_i \quad \mathcal{L}_{\Phi_i B_i \Sigma^*} = D_i \bar{\Sigma}^{*\mu} (\partial_\mu \Phi) B$$

Interaction kernels

$$V_{ji}^{\text{WT}} = -\frac{C_{ji}}{4f^2} (2\sqrt{s} - M_j - M_{4+5}), \quad V_{ji}^{\Sigma^*} = -\frac{D_j D_i}{3} \frac{|\mathbf{k}_j| |\mathbf{k}_i|}{\sqrt{s} - M_{\Sigma^*}}$$

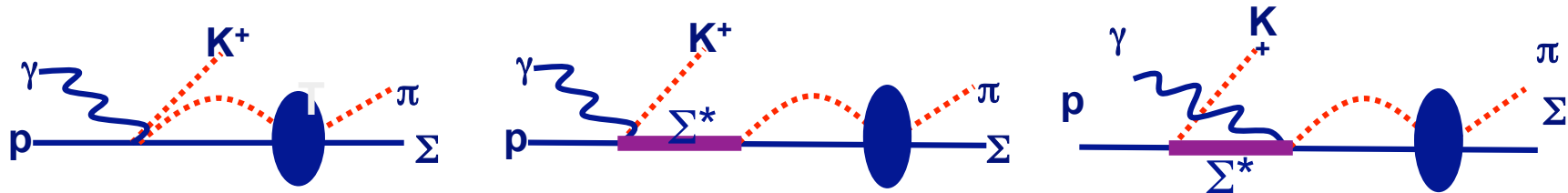
Unitarization

$$T_{ji} = T_{ii}^{\text{WT}} + [2(\hat{\mathbf{k}}_j \cdot \hat{\mathbf{k}}_i) - i\boldsymbol{\sigma} \cdot (\hat{\mathbf{k}}_j \times \hat{\mathbf{k}}_i)] T_{ii}^{\Sigma^*}$$

$$T^{\text{WT}} = \frac{1}{[1 - V^{\text{WT}} G]} V^{\text{WT}}, \quad T^{\Sigma^*} = \frac{1}{[1 - V^{\Sigma^*} G]} V^{\Sigma^*}$$

Theoretical analysis

$\Lambda(1405)$ photoproduction with (s,p)-wave contribution



Unitarized amplitude

$$T_\gamma = \underbrace{[i\boldsymbol{\sigma} \cdot (\hat{\boldsymbol{\epsilon}} \times \hat{\mathbf{k}}_\gamma)]}_{\text{s-wave}} T_a + \underbrace{[2(\hat{\mathbf{k}}_\pi \cdot \hat{\boldsymbol{\epsilon}}) - i\boldsymbol{\sigma} \cdot (\hat{\mathbf{k}}_\pi \times \hat{\boldsymbol{\epsilon}})]}_{\text{p-wave}} T_b + \underbrace{[2(\hat{\boldsymbol{\epsilon}} \cdot \hat{\mathbf{k}}_K) - i\boldsymbol{\sigma} \cdot (\hat{\boldsymbol{\epsilon}} \times \hat{\mathbf{k}}_K)]}_{\text{s-wave}} T_c$$

$\Phi B \rightarrow \Phi B$ total amplitudes

$$T_a = V_a + V_a G T_{WT},$$

$$T_b = V_b + V_b G T_{\Sigma^*},$$

$$T_c = V_c + V_c G T_{WT}.$$

$\Phi B \rightarrow \Phi B$ element amplitudes

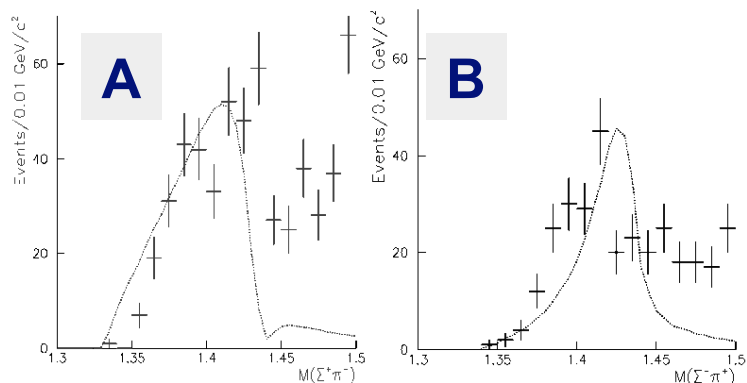
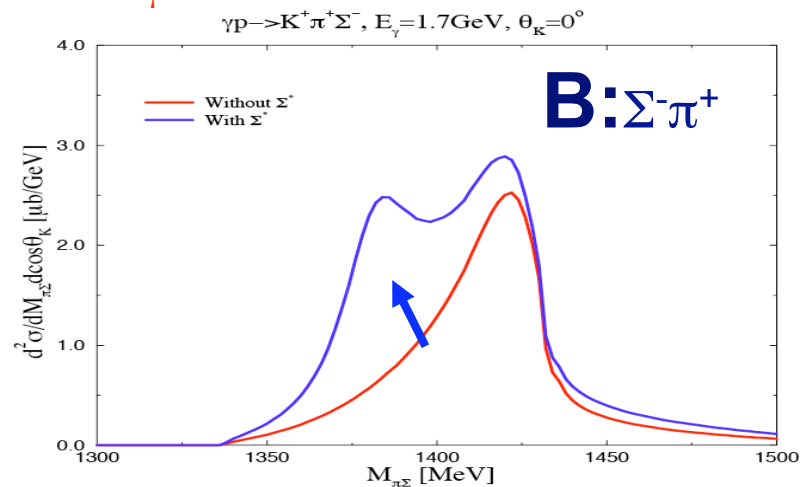
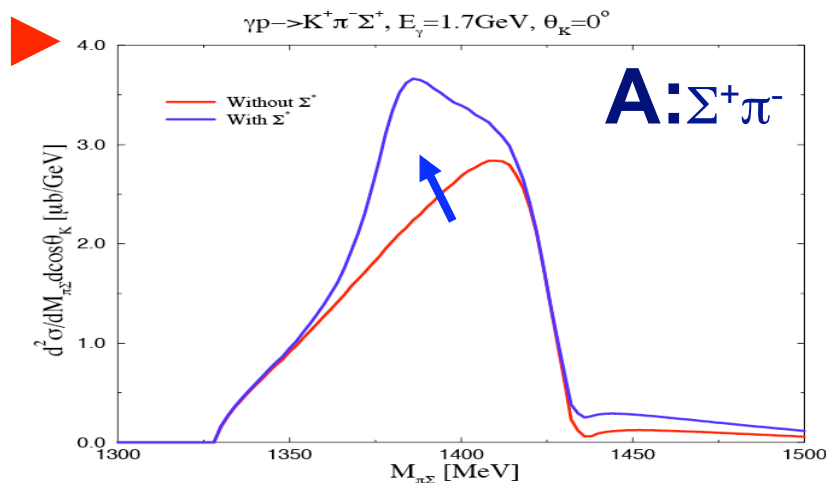
$$V_a = \frac{(e_\pi - e_K) C_{\bar{K}p \rightarrow \pi \Sigma} |\mathbf{k}_\gamma|}{8f^2 M_N},$$

$$V_b = \frac{e_K |\mathbf{k}_\pi|}{3} \frac{D_{\bar{K}p} D_{\pi \Sigma}}{E_\gamma + E_p - E_K - M_{\Sigma^*}},$$

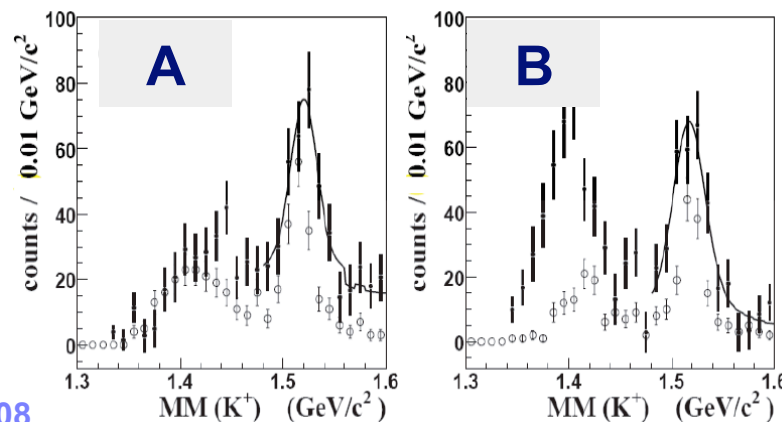
$$V_c = \frac{e_\pi |\mathbf{k}_K|}{3} \frac{D_{\bar{K}p} D_{\pi \Sigma}}{E_p - E_K - M_{\Sigma^*}},$$

Numerical results

$\Lambda(1405)$ photoproduction via $\gamma p \rightarrow K^+ \pi \Sigma$ at $E_\gamma = 1.7 \text{ GeV}$



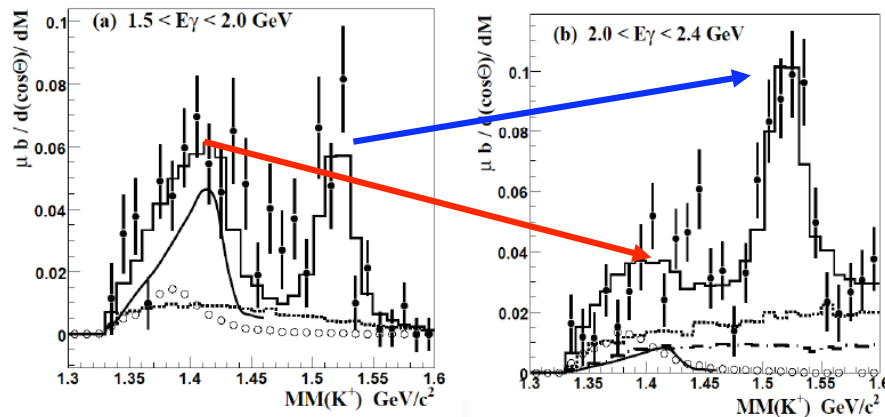
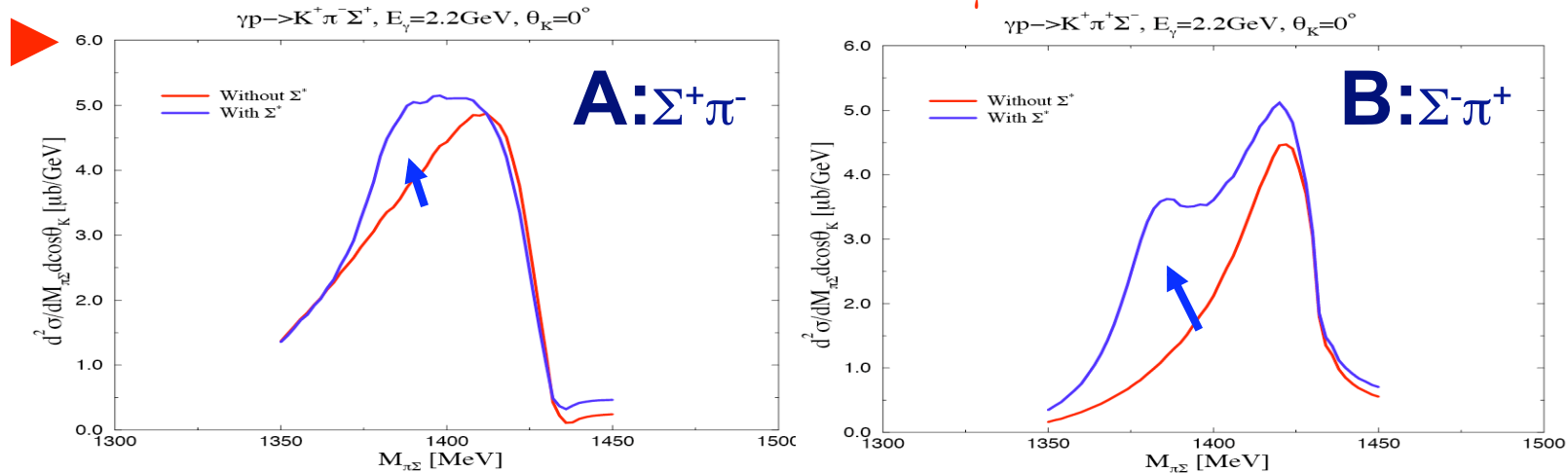
J.K.Ahn, Challenge to New Exotic Hadron at RCNP, 2008



Niiyama, PRC78,035202

Numerical results

$\Lambda(1405)$ photoproduction via $\gamma p \rightarrow K^+ \pi^- \Sigma$ at $E_\gamma = 2.2 \text{ GeV}$



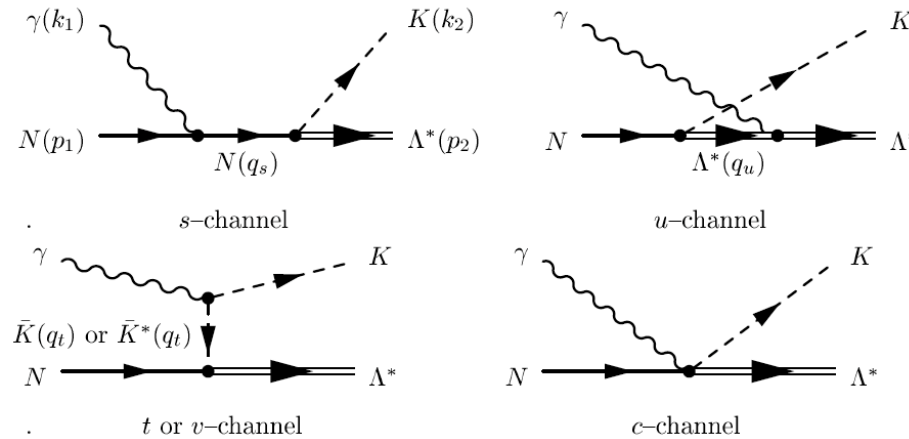
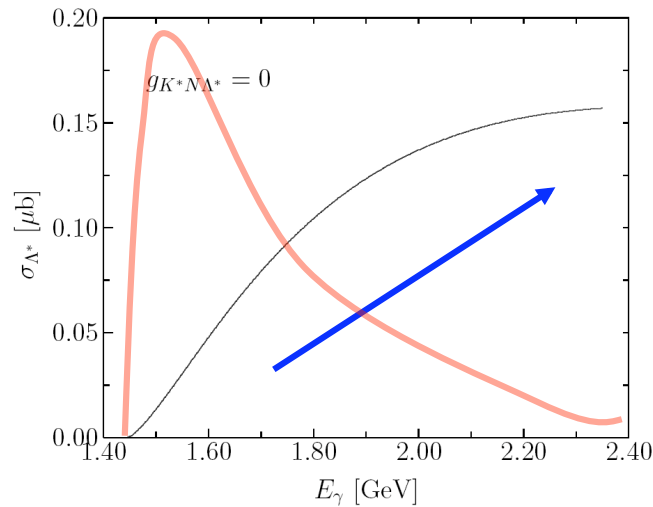
Niiyama, PRC78,035202

Strength of $M_{\pi\Sigma}$

- Theory: increasing
- Experiment: decreasing

Numerical results

$\Lambda(1405)$ photoproduction via $\gamma p \rightarrow K^+ \Lambda(1405)$ SiN et al, arXiv:0806.4029 [hep-ph]



Usual approaches can not reproduce the data

Indication of a strong resonance near the threshold?

cf) KKN molecular state D.Jido et al, arXiv:0806.3601 [nucl-th]

Numerical results

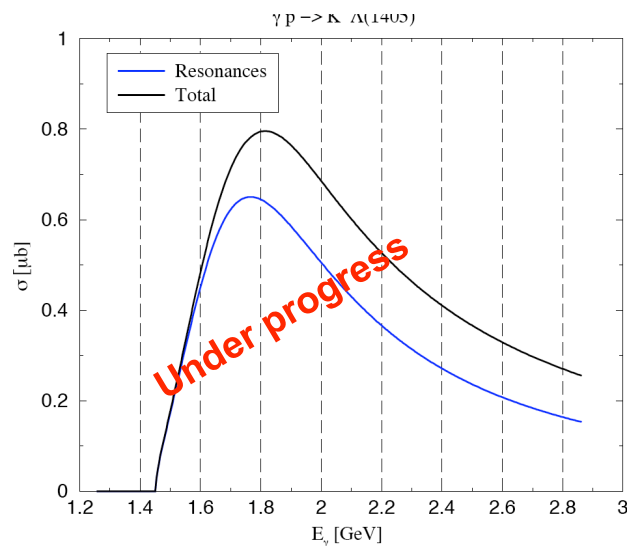
$\Lambda(1405)$ photoproduction via $\gamma p \rightarrow K^+ \Lambda(1405)$ SiN et al, arXiv:0806.4029 [hep-ph]

Considering resonances and two-pole structure

$$|\Lambda(1405), \text{Phys}\rangle = a|\text{high-}\Lambda(1429)\rangle + b|\text{low-}\Lambda(1392)\rangle + c|\text{uds}, I=0\rangle + \dots$$

$|\text{high-}\Lambda(1429)\rangle$ and $|\text{low-}\Lambda(1392)\rangle$: KN-bound state

$|\text{uds}, I=0\rangle$: NR-quark model (close)



	$ \Gamma_{N^* \rightarrow K\Lambda^*} $ [2]	$A_1^{P^*}$ [3]	$A_3^{P^*}$ [3]
$S_{11}(2030, 1/2^-)$	1.44 MeV	0.020 $\text{GeV}^{-1/2}$	—
$D_{13}(1960, 3/2^-)$	15.21 MeV	0.036 $\text{GeV}^{-1/2}$	-0.043 $\text{GeV}^{-1/2}$
$D_{13}(2055, 3/2^-)$	1.44 MeV	0.016 $\text{GeV}^{-1/2}$	0

Relative phase between the high and low ~ -1
(in ChUM by D.Jido)

Destructive interference between them!

Summary and perspective

$\Lambda(1405)$ photoproduction investigated to interpret recent LEPS with TPC

ChUM with s- and p-wave (Σ^* 1385) contributions

	Theory (ChUM)	Experiment
$\Sigma^-\pi^+$ vs. $\Sigma^+\pi^-$	$\Sigma^-\pi^+ \sim \Sigma^+\pi^-$	$\Sigma^-\pi^+ > \Sigma^+\pi^-$: Niiyama ($\Sigma^-\pi^+ \sim \Sigma^+\pi^-$: J.K.Ahn)
Energy dependence	Increasing	Decreasing
$\Sigma^*(1385)$ contribution	Strong	Weak

Possible explanations

- Other p-wave contributions: ground-state Y, M-B-loop contributions
- Larger K^+ angle: $0 < \cos\theta_K < 1$
- Resonance contributions & destructive interference of two-pole cont.
- $\Lambda(1405)$: mixture of various states: experimental evidence?
- How about $\Lambda(1520)$??

**Thank you very much
for your attention!!**

**Special thanks to
T.Nakano, J.K.Anh, M.Niiyama, and H.Kohri**