

# Hypernuclear weak decay :

-- present and future problems --

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# 1 Experiment : Pi-Mesonic Weak Decay

- Mesonic decay rates

Data ( ~ 2000 )

Hypernucleus		Ref.
${}^4_{\Lambda}\text{H}$	$\Gamma_{\pi^-}$	Outa(1995)
${}^4_{\Lambda}\text{He}$	$\Gamma_{\pi^0}, \Gamma_{\pi^-}$	Outa(1995,1998) Zeps(1998)
${}^5_{\Lambda}\text{He}$	$\Gamma_{\pi^0}, \Gamma_{\pi^-}$	Szymanski(1991)
${}^9_{\Lambda}\text{Be}$	$\Gamma_{\pi}$	Bando(1987) ?
${}^{11}_{\Lambda}\text{B}$	$\Gamma_{\pi}$ $\Gamma_{\pi^0}$ $\Gamma_{\pi^0}, \Gamma_{\pi^-}$	Grace(1985) Sakaguchi(1991) Noumi(1995)
${}^{12}_{\Lambda}\text{C}$	$\Gamma_{\pi^0}, \Gamma_{\pi^-}$ $\Gamma_{\pi^0}$ $\Gamma_{\pi^0}, \Gamma_{\pi^-}$	Szymanski(1991) Sakaguchi(1991) Noumi(1995)

## Data (2000~ )

Hypernucleus		Ref.
${}^{11}_{\Lambda}\text{B}$	$\Gamma_{\pi^{-}}$	Sato(2005)
${}^{12}_{\Lambda}\text{C}$	$\Gamma_{\pi^{-}}$	Sato(2005)
${}^{28}_{\Lambda}\text{Si}, {}^{27}_{\Lambda}\text{Si}$	$\Gamma_{\pi^{-}}$	Sato(2005)
${}_{\Lambda}\text{Fe}$	$\Gamma_{\pi^{-}}$	Sato(2005)
${}^7_{\Lambda}\text{Li}$	$\pi^{-}$ spectra, $\Gamma_{\pi^{-}}$	Botta(2008)
${}^9_{\Lambda}\text{Be}$	$\pi^{-}$ spectra, $\Gamma_{\pi^{-}}$	Botta(2008)
${}^{11}_{\Lambda}\text{B}$	$\pi^{-}$ spectra, $\Gamma_{\pi^{-}}$	Botta(2008)
${}^{15}_{\Lambda}\text{N}$	$\pi^{-}$ spectra, $\Gamma_{\pi^{-}}$	Botta(2008)

## 2 Experiment : Nonmesonic Weak Decay

- Decay rates,  $\Gamma_n/\Gamma_p$ ,  $\tau_{1/2}$   
Data

Hypernucleus		Ref.
${}^4_{\Lambda}\text{H}$	$\Gamma_{\text{nm}}$	Szymanski(1991) Outa(1998)
${}^4_{\Lambda}\text{He}$	$\Gamma_p, \Gamma_n$	Szymanski(1991) Outa(1998) Zeps(1998)
${}^5_{\Lambda}\text{He}$	$\Gamma_p, \Gamma_n$	Szymanski(1991) Noumi(1995)
${}^{11}_{\Lambda}\text{B}$	$\Gamma_{\text{nm}}, \Gamma_n/\Gamma_p$ $\Gamma_{\text{nm}}$	Szymanski(1991)Noumi(1995) Sato(2005)
${}^{12}_{\Lambda}\text{C}$	$\Gamma_{\text{nm}}, \Gamma_n/\Gamma_p$ $\tau_{1/2}$	Szymanski(1991)Noumi(1995) Sato(2005) Bhang(1998)Park(2000)

Hypernucleus		Ref.
${}^{27}_{\Lambda}\text{Al}$	$\Gamma_{\text{nm}}, \Gamma_{\text{n}} / \Gamma_{\text{p}}$	Sato(2005)
${}^{28}_{\Lambda}\text{Si}$	$\Gamma_{\text{nm}}, \Gamma_{\text{n}} / \Gamma_{\text{p}}$ $\tau_{1/2}$	Sato(2005) Bhang(1998) Park(2000)
${}_{\Lambda}\text{Fe}$	$\Gamma_{\text{nm}}, \Gamma_{\text{n}} / \Gamma_{\text{p}}$ $\tau_{1/2}$	Sato(2005) Bhang(1998) Park(2000)
${}_{\Lambda}\text{Bi}$	$\tau_{1/2}$	Kulesa(1998)

Data : neutron, proton Energy Spectra ( $N_n, N_p$ )  
n-n and n-p coincidence measurement,  $\Gamma_n/\Gamma_p$

Hypernucleus		Ref.
${}^5_{\Lambda}\text{He}$	$N_n, N_p$ n-n, n-p	Okada(2004) Outa(2005) Kang(2006)
${}^{12}_{\Lambda}\text{C}$	$N_n, N_p$ $N_p$ n-n, n-p	Okada(2004) Kim(2003) Hashimoto(2002) Outa(2005) Kim(2006)
${}^{89}_{\Lambda}\text{Y}$	$N_n$	Okada(2004) Kim(2003)

### 3 Theory : Mesonic weak decay

- $\pi$ -decay hamiltonian :

$$\mathbf{H}_\pi = \left[ \mathbf{s}_\pi \mathbf{X}^{(s)}(\mathbf{r}) + i \mathbf{p}_\pi \mathbf{X}^{(p)}(\mathbf{r}) \frac{(\vec{\sigma} \vec{\nabla})}{q_0} \right] \chi^{(-)*}(\vec{q}, \vec{r})$$

$$\mathbf{s}_{\pi 0} = -\mathbf{s}_{\pi^-} / \sqrt{2}$$

$$\mathbf{p}_{\pi 0} = -\mathbf{p}_{\pi^-} / \sqrt{2} \quad (\Delta I = 1/2 \text{ rule})$$

$$\mathbf{s}_{\pi^-} / \mathbf{p}_{\pi^-} = -3. \quad (\text{from } \pi \text{ decay asymmetry})$$

$$\vec{\Lambda} \rightarrow \mathbf{p} + \pi^-$$

$$\alpha_1^\pi = -0.64 \quad )$$

$$\chi^{(-)*}(\vec{q}, \vec{r}) : \pi\text{-on distorted wave}$$



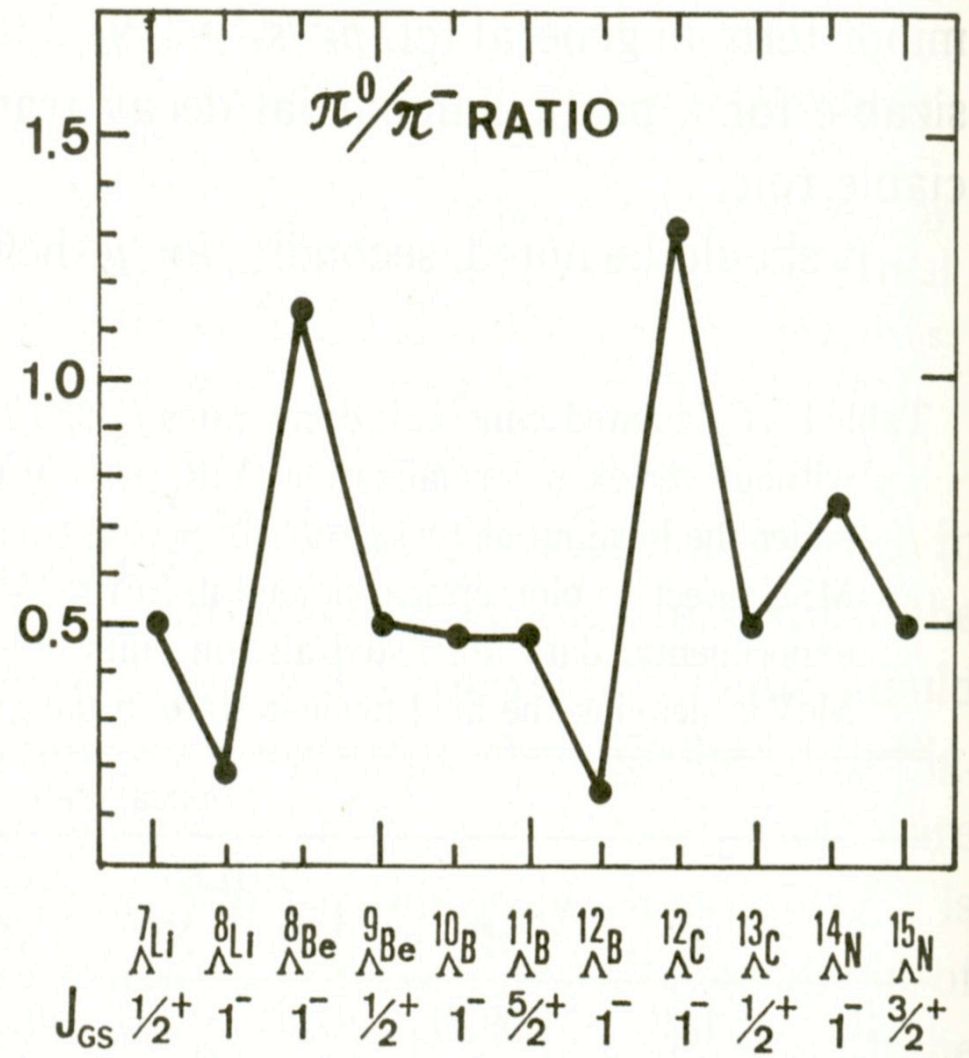
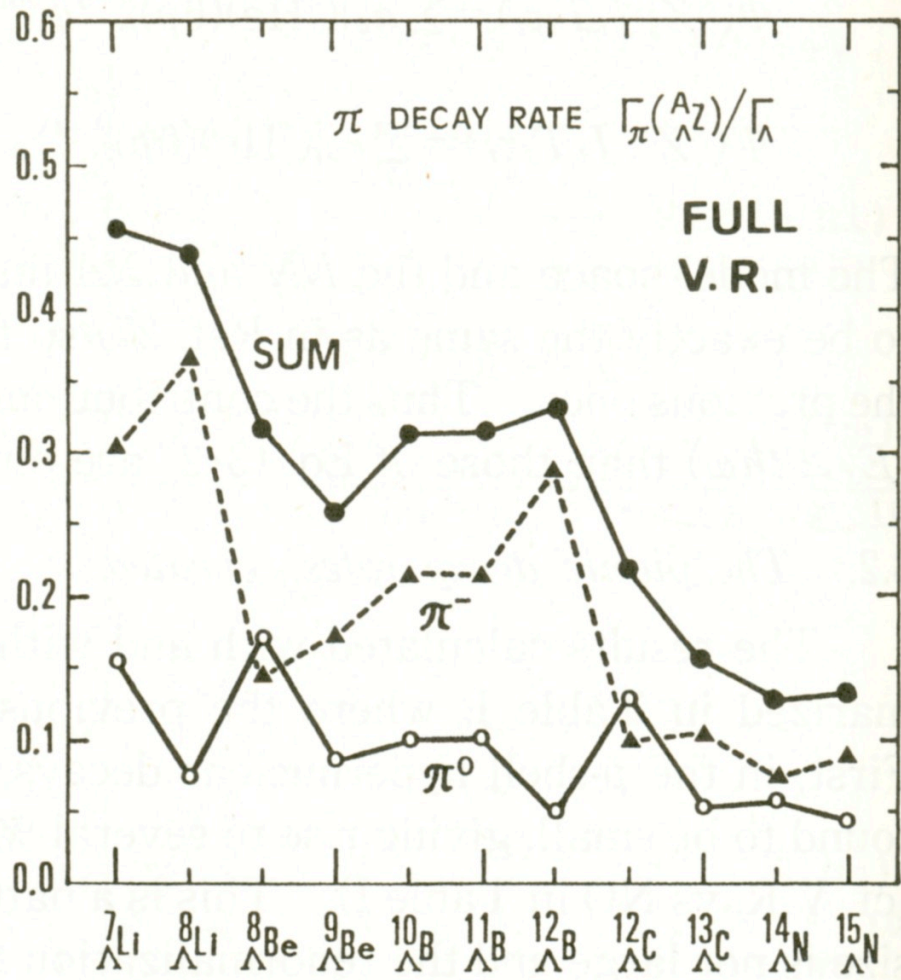
Vertex form factor (depends on  
 $\pi$ -Optical pot. )

$$\mathbf{X}^{(s)}(\mathbf{r}) = 1.0$$

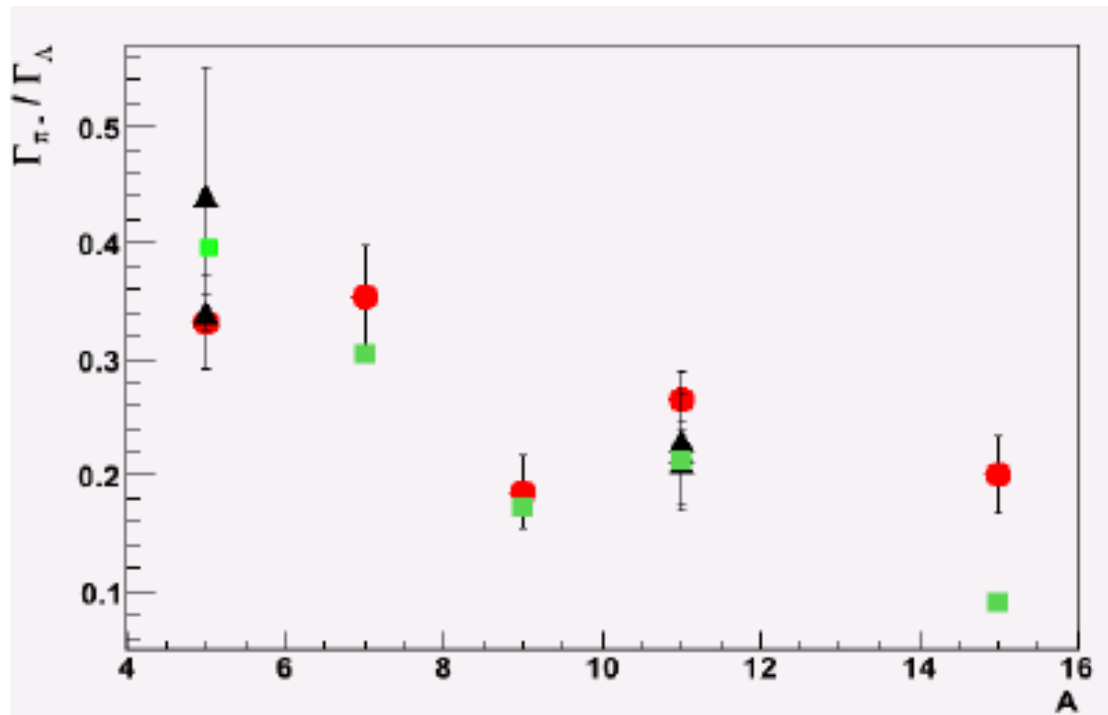
$$\mathbf{X}^{(p)}(\mathbf{r}) = L(\mathbf{r}) \quad (\text{LLEE for MSU pot.})$$

**Lorentz – Lorenz – Ericson – Ericson**

- $\Gamma_{\pi^0}, \Gamma_{\pi^-}, \pi^0/\pi^-, \Gamma_{\pi}^{\text{sum}}$
- Decay rates : sensitively depends on Hyp-structure
- Theor. Cal.  $\rightarrow$  good agreement with available data
- $A=4 : {}^4_{\Lambda}\text{H}, {}^4_{\Lambda}\text{He} \rightarrow \Gamma_{\pi}$ 's data favor Isle-type  
for  $V_{\Lambda\text{-nucleus}}$  pot.



- $A=15 : {}^{15}_{\Lambda}N$   $\pi^{-}$ -spectrum,  $\Gamma_{\pi^{-}}$  data  
(FINUDA 2008)  $\rightarrow$  cal. factor 2 small ?



- $\pi$ -on asymmetry from polarized hypernuclei

$${}^{\Lambda}_{\Lambda}Z(\mathbf{J}_i \mathbf{M}_i \mathbf{T}_i \tau_i, \mathbf{P}_H) \rightarrow {}^{\Lambda}Z'(\mathbf{J}_f \mathbf{T}_f \tau_f) + \pi(\vec{q})$$

- Angular distribution

$$\begin{aligned} & \Gamma_{\pi}(\mathbf{J}_i \mathbf{M}_i \mathbf{T}_i \tau_i, \mathbf{J}_f \mathbf{T}_f \tau_f ; \hat{q}) \\ &= \sum_{\mathbf{Q}} \Gamma_{\pi}^{\mathbf{Q}}(\mathbf{J}_i \mathbf{M}_i \mathbf{T}_i \tau_i, \mathbf{J}_f \mathbf{T}_f \tau_f ; \mathbf{q}) \mathbf{P}_{\mathbf{Q}}(\hat{q}) \\ &= \Gamma_{\pi}^{\mathbf{Q}=0}(\mathbf{J}_i \mathbf{M}_i \mathbf{T}_i \tau_i, \mathbf{J}_f \mathbf{T}_f \tau_f ; \mathbf{q}) (1 + \alpha_1 \mathbf{P}_1(\hat{q}) + \alpha_2 \mathbf{P}(\hat{q}) + \dots) \end{aligned}$$

$$\alpha_{\mathbf{Q}} = \frac{\Gamma_{\pi}^{\mathbf{Q}}}{\Gamma_{\pi}^{\mathbf{Q}=0}} \quad (\mathbf{M}_i = \mathbf{J}_i \text{ 100\% polarized})$$

$\alpha_1$  : asymmetry parameter

## ■ $\pi$ -on asymmetry

- ${}^5_{\Lambda}\text{He}$  : measured, Ajimura(1998)

Asymmetry  $A^{\pi} = P\alpha_1^{\pi}\varepsilon$

Measured  $A^{\pi}$  ,  $\varepsilon$  : reduct.factor (=0.81)

assumed  $\alpha_1^{\pi} = -0.64$  (free  $\Lambda$  val.)

→ Deduced Polarization P

$$P = 0.249 \text{ +/- } (\theta=2-7 \text{ deg})$$

$$P = 0.393 \text{ +/- } (\theta=7-15 \text{ deg}) < - > \text{ Consistent with cal.}$$

Theory (prediction)

NP. A489(1988)683

$${}^9_{\Lambda}\mathbf{B}(1/2^+) , \alpha_1^{\pi} = -0.64 \text{ (cal.)}$$

$${}^{13}_{\Lambda}\mathbf{C}(1/2^+) , \alpha_1^{\pi} = -0.16 \text{ (cal.)}$$

■  $\Delta I = 1/2$  rule for  $\Lambda \rightarrow N + \pi$  decay

- $\Lambda \diamond p + \pi^- \quad (63.9 \pm 0.5) \%$   
 $\quad \diamond n + \pi^0 \quad (35.8 \pm 0.5) \%$

Rule : established empirically

$$\frac{\Gamma_{\pi}^{\text{exp}}(\Lambda \rightarrow p + \pi^-)}{\Gamma_{\pi}^{\text{exp}}(\Lambda \rightarrow n + \pi^0)} \cong \frac{2}{1} = \frac{(1/2 \ 1/2 \ 1 - 1 | 1/2 - 1/2)^2}{(1/2 - 1/2 \ 1 0 | 1/2 - 1/2)^2}$$

\* Theoretical foundation, however,

**not yet clarified well**

\* see, Hiyama et al. PTP 112(2004)99

quark-quark correlations,  $(us)^0 \rightarrow (ud)^0$  considered.

\* also, Oka's group and other QCD works

## 4 Theory : $\Lambda N \rightarrow NN$ nonmesonic weak decay

- $\Delta S = 1$ ,  $Q = 176$ . MeV,  $q \sim 400$  MeV/c (free  $\Lambda$ )
- High momentum transfer process
- Short-range part of interactions contributes
- NMWD – dominant decay mode in medium-to-heavy  $\Lambda$ -hypernuclei

### A. Nonmesonic decay interactions, $V(\Lambda N - NN)$

Models :

1 one-pion exchange,  $V_\pi$

basic but not dominant

lightest  $0^-$  meson  $\rightarrow$  long-ranged, strong tensor

$\rightarrow$  fail to explain  $\Gamma_n/\Gamma_p$  (n/p –ratio) data

2 octet-meson exchanges :

$0^-$  pseudo-scalar ( $\pi, K, \eta$ ) exch.

$1^-$  vector ( $\rho, K^*, \omega$ ) exch.

$\pi + K$  : work additively for  ${}^3S_1 \rightarrow {}^3P_1$  ,

destructively for parity-cons. channels

→ enhance the n/p ratios, which explains features of exp. data (Good ! )

→ but not enough to explain  $\Gamma_{nm}$

other mesons : necessary to explain  $\Gamma_{nm}$

octet-meson exch. : not successful to explain asymmetry

parameter  $\alpha_\Lambda$  of  ${}^5_\Lambda\text{He}$  ,  ${}^{12}_\Lambda\text{C}$



3 correlated- $2\pi$  , uncorrelated- $2\pi$  exchanges :

$2\pi/\sigma$  :  $0^+$  scalar exch.

enhance the decay rates

$2\pi/\rho$  :  $1^-$  vector exch.

tensor force, opposite sign to  $V_\pi$

correlated- $2\pi$  + uncorrelated- $2\pi$  (+ octet-mesons) :

→ work favorably to explain  $\alpha_\Lambda$

( Chumillas et al. 2007)

4 Axial vector meson exchange :

$a_1$  :  $J^\pi = 1^+$ ,  $\leftrightarrow$  chiral partner of  $\rho$ , like  $\pi \leftrightarrow \sigma$

modeled as  $\rho\pi/a_1$  ,  $\sigma\pi/a_1$  -exch.

(+  $\pi$ ,  $K$ ,  $\omega$ ,  $2\pi/\sigma$ ,  $2\pi/\rho$ ) → work favorably for  $\alpha_\Lambda$

${}^5_\Lambda \mathbf{He}$  ,  ${}^{12}_\Lambda \mathbf{C}$

5 Direct quark interaction :

short-ranged

$\Delta I = 1/2$  &  $3/2$  contributions

$\Delta I = 3/2$  contributions, large for  $J = 0$  trans.

Direct quark int. alone  $\rightarrow$  not enough to explain  $\Gamma_{nm}$

Direct quark +  $\pi$  +  $K$  +  $\sigma$  :

$\rightarrow$  can explain  $\alpha_\Lambda$  of  ${}^5_\Lambda\text{He}$  (Sasaki et al. 2005)

6 Effective field theory :

low order effective field theory (LO pc +pv)

high mom. (short-distance) modes  $\rightarrow$  contact operator

$\pi, K \rightarrow$  treated as dynamical field, long-range part

$\rightarrow$  stress, importance of scalar-isoscalar contact int.

to fit data including  $\alpha_\Lambda$

${}^5_\Lambda\text{He}$  ,  ${}^{12}_\Lambda\text{C}$

## B. $\Gamma_n / \Gamma_p$ ( n/p ratio )

Exp.	${}^5_{\Lambda}\text{He}$	0.45 +/- 0.11 +/- 0.03	Kang(2006)
	${}^{12}_{\Lambda}\text{C}$	0.56 +/- 0.12 +/- 0.04	Outa(2005)
		0.51 +/- 0.13 +/- 0.05	Kim (2006)

Theory :

$V_{\pi} + V_K$  : important role to explain the large n/p ratios

mechanisms:

${}^3S_1 \rightarrow {}^3P_1$  , PV-channel (  $l = 1$  )

interference works additively

${}^{1,3}S \rightarrow {}^{1,3}S$  ,  ${}^3S_1 \rightarrow {}^3D_1$  , PC-channels

interferences work destructively

(  $V_{\pi} + V_K$  alone :  $\rightarrow$  not enough to explain  $\Gamma_{nm}$  )

## C. Asymmetry parameter $\alpha_1, \alpha_\Lambda$

Exp.	${}^5_\Lambda\text{He}$	0.11 +/- 0.08 +/- 0.04	Outa, Maruta(2005)
		0.07 +/- 0.08 (+0.08/-0.00)	Maruta(2006)
		0.24 +/- 0.22	Ajimura(2000)
	${}^{12}_\Lambda\text{C}$	-0.20 +/- 0.26 +/- 0.04	Outa, Maruta(2005)
		-0.16 +/- 0.28(+0.18/-0.00)	Maruta(2006)

Theory :

1 effective field theory (Parreno, 2004,2005)

$\alpha_\Lambda ({}^5_\Lambda\text{He})$  , fitted to data Ajimura(2000)

$A_1({}^{11}_\Lambda\text{B}) > 0$  , exp. - 0.20 +/- 0.10

$A_1({}^{12}_\Lambda\text{C}) > 0$  , exp. - 0.01 +/- 0.10

stressed : importance of scalar-isoscalar  
contact (short-ranged) interaction

## 2 $\sigma$ -exchange

Sasaki et al. (2005)

DQ +  $\pi$  + K +  $\sigma$ -exch.  $\rightarrow$  explain  $\alpha_\Lambda$  ( ${}^5_\Lambda\text{He}$ )

Barbero et al. (2006)

octet-meson +  $\sigma$ -exch.  $\rightarrow$  not succeed to explain  $\alpha_\Lambda$

Itonaga et al.  $\pi + 2\pi/\sigma + 2\pi/\rho + \omega + K$

$\rightarrow$  cannot explain  $\alpha_\Lambda$

$\rightarrow$  still controversial on  $\sigma$ -exchange

## 3 $2\pi$ -exchange (correlated & uncorrelated)

Chumillas et al. (2007) :

adopted  $V_{2\pi}$  by Jido et al.(2000)

Chumillas et al. (2007)

octet-meson + correl.- $2\pi$  ( $2\pi/\sigma$ ) + uncorrel.- $2\pi$

→ well explain  $\alpha_\Lambda$  of  ${}^5_\Lambda\text{He}$  &  ${}^{12}_\Lambda\text{C}$

Itonaga et al.  $\pi + 2\pi/\sigma + 2\pi/\rho + \omega + K$

→ cannot explain  $\alpha_\Lambda$

→ still controversial

4 axial vector  $a_1$ -exchange

$a_1$  : mass = 1230. MeV,  $J^\pi = 1^+$

$a_1(1^+) \leftrightarrow \rho(1^-)$ ,  $\pi(0^-) \leftrightarrow \sigma(0^+)$

■ new approach ( ? ), chiral partner mesons exch.

→ compatible with available data of  $\alpha_\Lambda$

## D. Problems (theory)

1 ●  $\alpha_\Lambda$  expression differs by authors? ( free  $\Lambda^+ p \rightarrow n + p$  )

W. A. Alberico, A.Ramos et al. (2005)

$$\alpha_\Lambda = \frac{2\sqrt{3} \operatorname{Re} [ ae^* - b(c - \sqrt{2}d)^* / \sqrt{3} + f(\sqrt{2}c + d)^* ]}{|a|^2 + |b|^2 + 3[|c|^2 + |d|^2 + |e|^2 + |f|^2]}$$

Sasaki, Izaki and Oka (2005)

$$\alpha_\Lambda = \frac{2\sqrt{3} \operatorname{Re} [ -ae^* - b(c - \sqrt{2}d)^* / \sqrt{3} + f(\sqrt{2}c + d)^* ]}{|a|^2 + |b|^2 + 3[|c|^2 + |d|^2 + |e|^2 + |f|^2]}$$

Itonaga et al.

$$\alpha_\Lambda = \frac{2\sqrt{3} \operatorname{Re} [ -ae^* + b(c - \sqrt{2}d)^* / \sqrt{3} + f(\sqrt{2}c + d)^* ]}{|a|^2 + |b|^2 + 3[|c|^2 + |d|^2 + |e|^2 + |f|^2]}$$

What is the origin of the difference?

2 ● necessary to check the proposed mechanism (model) to explain  $\alpha_\Lambda$

- contact int.
- $\sigma$ -exch.  $2\pi/\sigma$ , uncorrelated- $2\pi$
- $a_1$ -exch.
- direct-quark

3 ● New and different approach, possible ?

- ??
- strange-meson  $K_1(1400)$ ,  $J = 1^+$  ?

$\leftarrow \rightarrow K_1 = (\pi K^*) ?$

- role of  $\Delta I = 3/2$  ?
- What else ?

4 ● nonmesonic decay of  ${}_{\Lambda\Lambda}Z$  hypernuclei



# 5 J-PARC 実験に期待する

## A. Mesonic Decay

### 1 $\Gamma_{\pi^0}$ measurement of hypernuclei

:  ${}^7_{\Lambda}\mathbf{Li}$ ,  ${}^9_{\Lambda}\mathbf{Be}$ ,  ${}^{11}_{\Lambda}\mathbf{B}$ ,  ${}^{12}_{\Lambda}\mathbf{C}$ ,  ${}^{15}_{\Lambda}\mathbf{N}$ ,  ${}^{27,28}_{\Lambda}\mathbf{Si}$ ,  ${}^{56}_{\Lambda}\mathbf{Fe}$

\*  $\Gamma_{\pi^-}$ , measured at FINUDA & Sato et al.(2005)

\*  $\Gamma_{\pi^0}$  data of  ${}^{12}_{\Lambda}\mathbf{C}$   $\rightarrow$  still large error-bar

\* high quality data  $\leftrightarrow$  more informations for pion behavior or  $U_{\pi}^{\text{opt}}$  inside the nucleus

$$* \left| \langle \phi_{\mathbf{n}l\mathbf{j}}^{\mathbf{N}} \mid \tilde{\mathbf{j}}^{\pi} \mid \varphi_{0s}^{\Lambda} \rangle \right|^2$$

2  $\Gamma_{\pi^-}, \Gamma_{\pi^0}$  measurement of neutron-rich Hyp-nucl.

:  ${}^9_{\Lambda}\mathbf{He}$ ,  ${}^{10}_{\Lambda}\mathbf{Li}$

\* E-10 proposal (Sakaguchi)

\*  $\pi^-$  ( $\pi^0$ ) spectra may serve to determine the hypernuclear spin  $J^\pi$

3 Measurement of decay asymmetry  $\alpha_1^\pi$  of

:  ${}^9_{\Lambda}\mathbf{Be}$ ,  ${}^{13}_{\Lambda}\mathbf{C}$

\* Theoretical prediction exists for some typical Hyp.

\* weak decay mechanism and pion behaviors are rather well known

## B. Nonmesonic Decay

1 Measurement of decay rates,  $\Gamma_{nm}$ ,  $\Gamma_n/\Gamma_p$ , are desirable for p-shell and heavier Hy.

:  ${}^7_{\Lambda}\text{Li}$ ,  ${}^9_{\Lambda}\text{Be}$ ,  ${}^{10}_{\Lambda}\text{B}$ ,  ${}^{13}_{\Lambda}\text{C}$ ,  ${}^{16}_{\Lambda}\text{O}$ ,  ${}^{89}_{\Lambda}\text{Y}$ ,  ${}^{209}_{\Lambda}\text{Bi}$

- \* High quality data of n/p ratios exist only for  ${}^5_{\Lambda}\text{He}$  and  ${}^{12}_{\Lambda}\text{C}$
- \* mass-A dependence of  $\Gamma_{nm}$ ,  $\Gamma_n/\Gamma_p$  are known  $\rightarrow$  weak decay int. range will be deduced
- \* neutron-excess ( $N > Z$ ) effect on decay rates are studied

## 2 More asymmetry parameter $\alpha_{\Lambda}$ measurements are desirable

:  ${}^7_{\Lambda}\text{Li}(1/2^+)$ ,  ${}^{13}_{\Lambda}\text{C}(1/2^+)$ ,  ${}^{14}_{\Lambda}\text{N}(1^-)$

\*  $\alpha_{\Lambda}^{\text{NM}}$  's have HY-mass (shell) dependence or not ??

$\alpha_{\Lambda}^{\text{NM}}({}^5_{\Lambda}\text{He})$  and  $\alpha_{\Lambda}^{\text{NM}}({}^{12}_{\Lambda}\text{C})$  are sign different ?

\* What is the decisive mechanisms for the small  $\alpha_{\Lambda}^{\text{NM}}$  ?

-- What type of the decay interactions ?

-- final-state interactions ?

-- effect of  $\Delta I = 3/2$  ?

## 3 Measurement of $A = 4$ hypernuclei

:  ${}^4_{\Lambda}\text{H}$ ,  ${}^4_{\Lambda}\text{He}$

## Test of $\Delta = 1/2$ rule

$$\Gamma_{nm}({}^4_{\Lambda}\mathbf{H}) \sim 3R_{n1} + R_{n0} + 2R_{p0}$$

$$\Gamma_{nm}({}^4_{\Lambda}\mathbf{He}) \sim 2R_{n0} + 3R_{p1} + R_{p0}$$

$$\frac{\Gamma_n({}^4_{\Lambda}\mathbf{He})}{\Gamma_p({}^4_{\Lambda}\mathbf{H})} = \frac{R_{n0}}{R_{p0}} = 2 \quad (\text{if } \Delta I = 1/2)$$

## 4 Measurement of decay rate of double- $\Lambda$ hypernuclei

\*  $\Lambda\Lambda \diamond n\Lambda, p\Sigma^-, n\Sigma^+$

## 5 Hope to explore flavor nuclei including $\Lambda_c$

