# The Nature of Dark Energy and its Implications for Particle Physics and Cosmology

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## I. Introduction

• Current cosmological observations suggest that the universe is accelerating today.



In particular, "dark energy" almost dominates the present universe. However, we have not understood it yet.

## Plan of this talk

I. Introduction

2. Current constraints and some implications

3. The nature of dark energy affects other aspects of cosmology?

4. Summary

## 2. Current constraints and some implications

• What is dark energy?

Some "energy" which accelerates the present universe

**Cosmic acceleration:**  $\frac{\ddot{a}}{a} = -\frac{1}{6M_{\rm pl}^2}(\rho + 3P)$  $\begin{array}{c} a: {\rm scale factor} \\ \rho: {\rm energy \ density} \\ P: {\rm pressure} \end{array}$ 

If the equation of state  $w \equiv \frac{P}{\rho} < -\frac{1}{3}$ , the universe can be accelerated.

(However,  $w_m = 0$  for matter and  $w_r = 1/3$  for radiation)

 $\rightarrow$  We need something weird with  $w_x < -1/3$ .  $\rightarrow$  Dark Energy

#### Parameters characterizing dark energy

• Equation of state  $w_X = \frac{p_X}{\rho_X}$  (for a constant equation of state) • Energy density  $\Omega_X = \frac{\rho_X}{\rho_{\text{crit}}}$   $\rho_X(a) = \rho_{\text{crit}}\Omega_X a^{-3(1+w_X)}$ 

 $\star$  Dark energy can also fluctuate, thus we need more to characterize DE.

affects cosmic density fluctuation (CMB, LSS, ...)

- Speed of sound of DE  $c_s^2 \equiv \frac{\delta p}{\delta \rho}\Big|_{rest}$
- (• Anisotropic stress of DE  $\sigma_X$  )

 $\star$  These quantities can be constrained by observations.



(First we discuss the effects of the background modification.)

**\bigstar** Once  $\Omega_X$  and  $w_X$  are specified, we can know how the universe is accelerated.

$$\frac{\ddot{a}}{a} = -\frac{\rho_{\rm crit}}{6M_{\rm pl}^2} \left( \Omega_m a^{-3} + \Omega_r a^{-4} + \Omega_K a^{-2} + \frac{\Omega_X a^{-3(1+w_X)}}{\Omega_X a^{-3(1+w_X)}} \right)$$

(Or the background evolution)

$$H^{2}(a) = H_{0}^{2} \left( \Omega_{m} a^{-3} + \Omega_{r} a^{-4} + \Omega_{K} a^{-2} + \frac{\Omega_{X} a^{-3(1+w_{X})}}{\Omega_{X} a^{-3(1+w_{X})}} \right)$$

 $\bigstar$   $\Omega_X$  and  $w_X$  can be constrained by various observations such as

- Type la supernovae (SNela)
- Cosmic microwave background (CMB)
- Baryon acoustic oscillation (BAO)

• SNela

Luminosity distance 
$$d_L = \frac{1+z}{\sqrt{|\Omega_k|}} S\left(\sqrt{|\Omega_k|} \int_0^z \frac{dz'}{H(z')/H_0}\right)$$

[Recent observations: A. G. Riess et al., arXiv:astro-ph/0611572; Davis et al., arXiv:astro-ph/0701510; Wood-Vasey et al, astro-ph/0701041]

#### • CMB

Acoustic peak position gives the angular diameter distance to last scattering surface

Angular diameter distance  $d_A = \frac{1}{1+z} \int \frac{dz'}{H(z')}$ 

[Spergel et al., astro-ph/0603449; Wang&Mukherjee, astro-ph/0604051]

#### • BAO

Baryon acoustic peak was detected by observing galaxy samples  $z \sim 0.35$ .

[Eisenstein et al, Ap] 633, 560, 2005]







## Observational Constraints on the eq. of state

#### • For a constant equation of state (and a flat univ.)



# Constraint from the shift parameter (position of acoustic peak)



**<u>Recent constraints</u>** (for a flat universe and a constant equation of state.)

• Spergel et al (WMAP3), 2006

[WMAP3+other CMB+2dF+SDSS+SN]  $w_X = 0.926^{+0.054}_{-0.053}$  (with perturbation)

[WMAP3+BAO]  $w_X = -0.86^{+0.25}_{-0.23}$  (with perturbation)

[WMAP3+SN(gold)]  $w_X = -0.919^{+0.081}_{-0.080}$  (with perturbation)

[WMAP3+SN(SNLS)]

 $w_X = -0.967^{+0.073}_{-0.072}$  (with perturbation)

- Tegmark et al, PRD74, I 23507, 2006. (astro-ph/0608632) [WMAP3+SDSS]  $w_X = -1.00^{+0.17}_{-0.19}$  (with perturbation)
- Percival et al, arXiv:0705.3323 [astro-ph]
   [CMB(theta)+SN(SNLS)+BAO(SDSS+2dF)]

 $w_X = -1.004 \pm 0.089$ 

• Wood-Vasey et al, astro-ph/0701041 [SN(ESSENCE)+BAO(SDSS)]  $w_X = -1.05^{+0.13}_{-0.12}$ 

## Models for an accelerating universe proposed so far

- (No criterion for the choice of models here.)
- Cosmological constant (A) ( $w_x = -1$ )
- Scaler field
  - Quintessence
  - · K-essence (with a non-canonical kinetic term)
- · DGP (Dvali-Gabadaze-Porrati) model
- f(R) gravity
- Cardassian model ( $H^2 = \frac{1}{3M_{pl}^2}(\rho + B\rho^n)$ )
- Ghost condensate



## **★** Cosmological constant?

- A naive estimate of  $\Lambda$  in quantum field theory:  $\rho_{\Lambda} \sim M_{\rm pl}^4 \sim (10^{18} {\rm GeV})^4$
- From cosmological observations:  $\rho_{\Lambda} \sim \Lambda^4 \sim (10^{-3} {\rm eV})^4$

 $\longrightarrow M_{\rm pl}^4 / \Lambda^4 \sim 10^{120}$  !

## • Numerical coincidence?



## ★ Dark energy is dynamical?

## An example: quintessence (a scalar field) Q

• "Tracker type" model

(the ratio  $\rho_{\rm DE}/\rho_{\rm m}(\rho_{\rm rad})$  constant)

•  $V(Q) = \Lambda^{4+\alpha}/Q^{\alpha}$ 

[Ratra,Peebles PRD 37,3406,1998]

•  $V(Q) = \Lambda^4 \exp(-\lambda Q)$ 

[Ferreira,Joyce PRD 58,023503,1998]

- $V(Q) = V_0 \left[ (Q B)^2 + A \right] \exp(-\lambda Q)$ [Skordis, Albrecht PRD 66,043523,2002]
- $V(Q) = (\Lambda^4/Q^{\alpha}) \exp(\kappa Q^2)$

[Brax,Martin PLB 468,40,1999]

- "Tracker oscillating" model
  - $V(Q) = \Lambda^4 \left[1 + A\sin(\nu Q)\right] \exp(-\lambda Q)$

[Dodelson,Kaplinghat,Stewart PRL 85,5276,2000]

★ Various models even just for quintessence....

• "cosine type" model

(Psuedo-Nambu-Goldston boson)

$$V(Q) = \Lambda^4 \left(1 - \cos(Q/f_Q)\right)$$

[Coble,Dodelson,Frieman PRD 55,1851,1997; Viana,Liddle PRD 57,674,1998]

(mass scale: 
$$m_Q \sim H_0 \sim 10^{-33} {\rm eV}$$
 )



 $\star$  (Too) Many models proposed (but none of them are compelling).

Phenomenological approach using observations

## ★ Key questions:

• The equation of state is - I (cosmological constant) or not?

• The equation of state is time-dependent  $(dw_x/dz \neq 0)$  or not ?

• Probe the time-dependence of the equation of state  $W_x(z)$ .

Although time dependence of *W<sub>x</sub>* can be complicated, in most analysis, some simple parametrizations are adopted. (and see *dW<sub>x</sub>/dz*)

Observational Constraints on the eq. of state II

 $\star$  For the time-varying equation of state

• Assuming a simple model:

$$w_X = w_0 + w_1(1-a) = w_0 + w_1 \frac{z}{1+z}$$
 (A flat universe is assumed.)



Observational Constraints on the eq. of state III

• Another parametrization [Hannestad, Mortsell JCAP0409,011,2004]

$$w(a) = w_0 w_1 \frac{a^q + a_s^q}{a^q w_1 + a_s^q w_0}$$
 (LCDM:  $w_0 = w_1 = -1$ )



Observational Constraints on the eq. of state IV

• Another parametrization [Wetterich PLB 594, 17, 2004]

$$w_X(z) = \frac{w_0}{[1+b\log(1+z)]^2}$$
 (LCDM:  $w_0 = -1, b = 0$ )



[Movahed, Rahvar PRD 73, 083518, 2006]

#### (Short summary of this part)

 A cosmological constant is allowed (in almost all analysis). (looks pretty good in some analysis/data set.)

• Various dark energy models are still allowed.

 Many candidates for dark energy (accelerating universe) have been proposed, (fortunately/unfortunately) none of them are compelling.

(• For dark matter, plausible candidates in particle physics.)

★ In fact, more to be specified to characterize dark energy

• We also have to specify the perturbation nature of dark energy.

- Speed of sound of DE  $c_s^2 \equiv \frac{\delta p_X}{\delta \rho_X}\Big|_{\text{rest}}$ .
- Anisotropic stress of DE  $\sigma_X$

[Hu ApJ 506, 485, 1998; Bean&Dore PRD 69,083503,2003]

[Hu ApJ 506, 485, 1998; Koivisto&Mota PRD 73, 083502, 2006; Ichiki&TT PRD, astro-ph/0703549]

$$\sigma'_X + 3\mathcal{H}\frac{c_a^2}{w_X}\sigma_X = \frac{8}{3}\alpha\left(\theta_X + \frac{h'}{2} + 3\eta'\right)$$
: specified by "viscosity" parameter  $\boldsymbol{\Omega}$ 

• Perturbation equations for <u>a general dark energy</u>

For density pert.,  $\delta'_X = -(1+w_X) \left[k^2 + 9\mathcal{H}^2(c_s^2 - c_a^2)\right] \frac{\theta_X}{k^2} - 3\mathcal{H}(c_s^2 - w_X)\delta_X - (1+w_X)\frac{h'}{2}$ 

For velocity pert.,  $\theta'_X = -\mathcal{H}(1 - 3c_s^2)\theta_X + \frac{c_s^2k^2}{1 + w_X}\delta_X - k^2\sigma_X$ 

(the prime denotes a derivative w.r.t. the conformal time)

where 
$$c_a^2 \equiv \frac{p'_X}{\rho'_X} = w_X - \frac{w'_X}{3\mathcal{H}(1+w_X)}$$
.  $\checkmark$  Specified once EoS is given

These can affect cosmic density density fluctuation

#### Effects of fluctuation of dark energy



• For quintessence (a scalar field with the canonical kinetic term),  $\,\,c_s^2=1$ 

- For example, k-essence (a scalar field with a non-canonical kinetic term),  $c_s^2 
  eq 1$
- $\star$  Low multipoles of CMB are mainly affected.

 $\star$  Even if the eq. of state is same, fluctuation make some difference.

 $\star$  Cross-correlation of ISW and LSS may be helpful.

Although the sound speed and anisotropic stress themselves cannot be severely constrained, constraints on other quantities can be affected.

• Constraint on w with different assumptions for  $\,c_s^2$  and  $\,lpha$  .



<sup>[</sup>Ichiki&TT PRD, astro-ph/0703549]

 $\star$  The assumption may cause ~10 % difference.

• Constraint on w with different assumptions for  $c_s^2$  and lpha .



[Ichiki&TT PRD, astro-ph/0703549]

## 3. The nature of dark energy affects other aspects of cosmology?

- ★ Main concern of dark energy research is to figure out what the dark energy is.
  - the nature of dark energy can affect other aspects?
- Implications for other aspects? (Some examples)
  - Dark energy and mass varying neutrinos
    - Next Takahashi's talk
  - Relic abundance of DM can be affected by quintessence model?
  - The curvature of the universe
  - Primordial fluctuation (spectral index, gravity wave, scalar spectral running)

(Example I)

## • <u>Relic abundance of DM in models with quintessence model</u>

In some models of quintessence, the kinetic
 energy of quintessence can dominate the universe.

[Salati PLB 571,121,2003; Rosati PLB 570,5,2003; Profumo & Ullio JCAP0311,006,2003; Pallis JCAP 0510,015,2005]

• During kinetic energy-dominated phase,  $ho_{\phi} \propto a^{-6}$ 

Example:  $V(\phi) = M_{\rm pl}^4 \exp\left(-\lambda \phi/M_{\rm pl}\right)$ 



# ★ Relic abundance can be enhanced when DM decouple in kination domination period.

- Standard case  $\rightarrow$  DM decouples during radiation-dominated epoch  $\left(H^2 = \frac{1}{3M_{rl}^2}\rho \propto a^{-4}\right)$
- If DM decouples during Kination domination (kinetic energy dominated epoch)

• 
$$H^2 = \frac{1}{3M_{\rm pl}^2} \rho \propto a^{-6}$$
 •  $\frac{dn_{\chi}}{dt} + 3Hn_{\chi} = \langle \sigma v \rangle \left\{ (n_{\chi}^{(0)})^2 - n_{\chi}^2 \right\}$ 



(Example 2)

#### Effects on the constraint on the curvature of the universe

• It is usually said that current observations favor a flat universe.



• The flatness is robust even if we assume different types of dark energy?

#### • The EoS of dark energy also affect the CMB power spectrum.



Degeneracy between EoS of DE and the curvature



• Assuming a cosmological constant as dark energy

[Ichikawa & TT, PRD 73, 083526 (2006)]

ullet Assuming a constant equation of state  $w_X$ 

( $w_X$  is marginalized over.)



Contours of EoS giving the minimum chi2 for each  $(\Omega_m, \Omega_X)$ .



Each observation favors different values of EoS.

When all data combined, it gives a flat universe.

• Assuming a time-varying equation of state as



$$w_X = w_0 + (1 - a)w_1$$
$$= w_0 + \frac{z}{1 + z}w_1$$

( $w_0$  and  $w_1$  are marginalized over.)

A flat universe is favored even though we assume a time-varying equation of state.

[Ichikawa & TT, PRD 73, 083526 (2006)]

• The flatness is robust?

• Assuming a time-varying equation of state as



[Ichikawa, Kawasaki, Sekiguchi & TT, JCAP 12, 005(2006)]

• Assuming a time-varying equation of state as



(Example 3)

• Effects on the constraints on primordial fluctuation

★ The nature of primordial fluctuation is now severely constrained by observations.

- Spectral index  $n_s: P_\mathcal{R} \propto k^{n_s-1}$
- Tensor-to-scalar ratio  $r: r = \frac{P_T}{P_P}$
- Running of the scalar spectral index lphas  $\,$  :  $lpha_s$  =

$$a_s = \frac{d\ln n_s}{d\ln k}$$



Running of scalar spectral index



#### • Effects of dark energy and primordial fluctuation look similar?



#### Dark energy fluctuation and the running of spectral index



#### Constraint on the scalar spectral index and tensor-to-scalar ratio

• For some different values of  $\,c_s^2$  and lpha

[Ichiki & TT in preparation]

(Using MCMC approach)



• The nature of DE can also affect the constraints on the running spectral index.

## <u>4. Summary</u>

- Current precise cosmological observations can constrain the equation of state for dark energy, severely in some cases.
- A cosmological constant is allowed (in almost all analysis), however, various kinds of dark energy can also still be allowed.
- The nature of dark energy also affects other aspects of cosmology (such as constraints on the curvature, primordial fluctuation, ....).
- Dark energy is one of the most important problems in today's science. We need to keep working on it.