# 研究会「宇宙初期における時空と物質の進化」

# 元素合成の現状

# Big Bang Nucleosynthesis (BBN)

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# I. Introduction BBN: Theory vs. Observation



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# I. Introduction BBN: Theory vs. Observation





- Introduction
- He4
- Li7
- Li6
- D

2. He4



# Energy Level Diagram of Hel



Benjamin, Skillman, Smits 1999, ApJ 514,307

# Spectrum

#### MRK 193 Izotov, Thuan, Lipovetsky (1994)



# Abundance of singly ionized Helium $y^+ = n(\text{HeII})/n(\text{HII})$ $y^{+} = \frac{F(\lambda)}{F(H\beta)} \frac{E(H\beta)}{E(\lambda)} \left(\frac{W(H\beta)}{W(H\beta) + a_{\rm HI}}\right) \left(\frac{W(\lambda) + a_{\rm HeI}}{W(\lambda)}\right) 10^{f(\lambda)C(H\beta)} \frac{1}{f_{\lambda}}$ • $\frac{E(H\beta)}{E(\lambda)}$ :Theoretical emissivity scaled to $H\beta$ • $F(\lambda)$ : observed line intensity $a_{\rm HI}, a_{\rm HeI}$ : underlying stellar absorption • $W(\lambda)$ : equivalent width $L_{H\beta} = W(H\beta)L_{\lambda}(\lambda 4861)$ • $f(\lambda)C(H\beta)$ : extinction relative to H $\beta$ • $f_{\lambda}(\tau)$ : optical depth function with collisional correction



Reddening and Stellar absorption

scattering and absorption Reddening (extinction) by interstellar dust  $I_{\lambda} = I_{\lambda 0} e^{-\tau_{\lambda}}$  $au_{\lambda} = Cf(\lambda)$  extinction law  $\log \left[ \frac{I(\lambda)}{I(H\beta)} \right] = \log \left| \frac{F(\lambda)}{F(H\beta)} \right| + C(H\beta)f(\lambda)$ intrinsic line observed line intensity intensity Underlying stellar absorption

Solving for reddening and underlying absorption

 $H\alpha/H\beta, H\gamma/H\beta, H\delta/H\beta \Rightarrow C(H\beta), a_{HI}$ 

correction for stellar absorption

 $F_A(\lambda) = F(\lambda) \left( \frac{W(\lambda) + a_{\rm HI}}{W(\lambda)} \right)$  W: EW (equivalent width)

 $\underline{L}_{H\beta} = W(H\beta)L_{\lambda}(\lambda 4861)$ 

reddening correction

 $X_R(\lambda) = \frac{I(\lambda)}{I(H\beta)} = \frac{F_A(\lambda)}{F_A(H\beta)} 10^{f(\lambda)C(H\beta)}$ 

theoretical value

 $X_T(6563) = 0.3862(\log T_4)^2 - 0.4817\log T_4 + 2.86\dots$ 

 $T_4 \equiv T/10^4 K$ 

• take 
$$\chi^2$$
 minimum  
 $\chi^2 = \sum_{\lambda} \frac{(X_R(\lambda) - X_T(\lambda))^2}{\sigma_{X_R}^2(\lambda)}$ 



# Abundance of singly ionized Helium $y^+ = n(\text{HeII})/n(\text{HII})$



# Theoretical emissivities

## Benjamin, Skillman, Smits 1999, ApJ 514,307 [BSS]

 $\overline{|E(H\beta)/E(3889)} = 0.904T^{-0.1}\overline{73-0.00054n_e}$ 

 $E(H\beta)/E(4026) = 4.297T^{0.090-0.000063n_e}$ 

 $E(H\beta)/E(4471) = 2.010T^{0.127} - 0.00041n_e$ 

 $E(H\beta)/E(5876) = 0.735T^{0.230-0.00063n_e}$ 

 $E(H\beta)/E(6678) = 2.580T^{0.249-0.00020n_e}$ 

 $E(H\beta)/E(3889) = 12.45T^{-0.917}$ 

 $/ [3.494 - (0.793 - 0.0015n_e + 0.000000696n_e^2)T]$ 

# Helium Abundance

$$y^{+} = \frac{F(\lambda)}{F(H\beta)} \frac{E(H\beta)}{E(\lambda)} \left(\frac{W(H\beta)}{W(H\beta) + a_{\rm HI}}\right) \left(\frac{W(\lambda) + a_{\rm HeI}}{W(\lambda)}\right) 10^{f(\lambda)C(H\beta)} \frac{1}{f_{\lambda}}$$

$$\bar{y} = \sum_{\lambda} \frac{y^{+}(\lambda)}{\sigma(\lambda)^{2}} / \sum_{\lambda} \frac{1}{\sigma(\lambda)^{2}} \quad \text{minimize } \chi^{2} \quad \bullet \quad \text{determine parameters} \\ (T), n_{e}, a_{\text{HeI}}, \tau \\ [T = T(\text{OIII})] \\ \chi^{2} = \sum_{\lambda} \frac{(y^{+}(\lambda) - \bar{y})^{2}}{\sigma(\lambda)^{2}} \quad \Delta\chi^{2} = 1 \quad \bullet \quad \text{uncetainties in parameters}$$

# Temp. measurement from [OIII] lines



# Spectrum

#### MRK 193 Izotov, Thuan, Lipovetsky (1994)



# **Recent Works**

- Izotov & Thuan 1998, 2004
  - 45 (89) low metallicity HII regions
  - use [OII] emission lines to determine T T(HeII) = T(OIII)

 $Y_p = 0.244 \pm 0.002$ 

- Peimbert, Peimbert & Ruitz 2000
  - HII region NGC 346 in SMC
  - use Hel emission line to determine T = T(

T(HeII) < T(OIII)

• Luridiana et al 2003

• 5 metal poor HII regions

$$Y_p = 0.239 \pm 0.002$$

 $Y_p = 0.2345 \pm 0.0026$ 

# Izotov, Thuan 2004



Fig. 2.— Linear regressions of the helium mass fraction Y vs. oxygen and nitrogen abundances for a total of 82 H II regions in 76 blue compact galaxies. In panels a) and b), Y was derived using the 3  $\lambda$ 4471,  $\lambda$ 5876 and  $\lambda$ 6678 He I lines, and in panels c) and d), Y was derived using the 5  $\lambda$ 3889,  $\lambda$ 4471,  $\lambda$ 5876,  $\lambda$ 6678 and  $\lambda$ 7065 He I lines.

	Number of	Oxygen		Nitrogen	
Method	H II Regions	Regression	$\sigma$	Regression	$\sigma$
3 He I lines <sup>a, b</sup>	2 45 89	$0.2451 \pm 0.0018 + 21 \pm 21(O/H)$ $0.2429 \pm 0.0009 \pm 51 \pm 9(O/H)$	0.0048	$0.2452 \pm 0.0012 + 603 \pm 372 (N/H)$ $0.2439 \pm 0.0008 \pm 1063 \pm 183 (N/H)$	0.0044
5 He I lines <sup>c,c</sup> 5 He I lines <sup>c,c</sup>	l $7l$ $7$	$0.2421\pm0.0003 + 51\pm 5(O/H)$ $0.2421\pm0.0021 + 68\pm22(O/H)$ $0.2444\pm0.0020 + 61\pm21(O/H)$	0.0040 0.0035 0.0040	$\begin{array}{r} 0.2435 \pm 0.0006 + 1003 \pm 103(N/H) \\ 0.2446 \pm 0.0016 + 1084 \pm 442(N/H) \\ 0.2466 \pm 0.0016 + 954 \pm 411(N/H) \end{array}$	0.0031 0.0040 0.0044

<sup>a</sup>Data are from IT98.

<sup>b</sup>Only collisional and fluorescent enhancements are taken into account. We have adopted  $T_e$ (He II) =  $T_e$ (O III) and ICF(He) = 1.

<sup>c</sup>Collisional and fluorescent enhancements of the He I lines, collisional excitation of hydrogen lines, underlying He I stellar absorption and differences between  $T_e$  (He II) and  $T_e$  (O III) are taken into account. ICF (He) is set to 1.

<sup>d</sup>Calculated with  $EW_a(H8 + He I 3889) = 3.0\text{\AA}, EW_a(He I 4471) = 0.4\text{\AA}, EW_a(He I 5876) = 0.3 EW_a(He I 4471), EW_a(He I 6678) = EW_a(He I 7065) = 0.1 EW_a(He I 4471).$ 

<sup>e</sup>Calculated with  $EW_a(H8 + He I 3889) = 3.0\text{\AA}$ ,  $EW_a(He I 4471) = 0.5\text{\AA}$ ,  $EW_a(He I 5876) = 0.3 EW_a(He I 4471)$ ,  $EW_a(He I 6678) = EW_a(He I 7065) = 0.1 EW_a(He I 4471)$ .

# $Y_p = 0.244 \pm 0.002$

# T(Hell)/T(OIII)

#### Peimbert, Peinbert, Luridiana (2002)



FIG. 1.—The ratio  $T_e(\text{He II})/T_e(\text{O III})$  as a function of  $T_e(\text{O III})$  and temperature fluctuations for the case in which all the O is O<sup>++</sup>. When O<sup>+</sup> is present, higher  $t^2$  values are expected, particularly for those objects with the highest  $T_e(\text{O III})$  values (see Fig. 2). Typical  $t^2$  values in H II regions are in the 0.01–0.04 range.

average temp  $T_0 = \frac{\int T n_e n_p dV}{\int n_e n_p dV}$ mean square temp variation  $t^2 = \frac{\int (T - T_0)^2 n_e n_p dV}{T_0^2 \int n_e n_p dV}$ pure OIII nebula

 $T(\text{HeII}) = T(\text{OIII}) \left[ 1 - \left( \frac{90800}{T(\text{OIII})} - 0.2 \right) \frac{t^2}{2} \right]$ 

# Recent Works (cont.)

- Olive & Skillman 2004
  - 7 HII regions of IT98
  - use Hel emission lines to determine T
  - underlying stellar absorption

 $Y_p = 0.249 \pm 0.009$ 

- Fukugita, MK 2006
  - 33 HII regions of IT04
  - use OIII emission line to determine T

underlying stellar absorption

 $Y_p = 0.250 \pm 0.004$ 

# Olive, Skillman 2004



# Helium Abundance in HII region



Fukugita,Kawasaki (2006)

# Without stellar absorption



Fukugita, Kawasaki (2006)

 $- Y_p = 0.234 \pm 0.004$ 



# New Determination of Y<sub>P</sub> Use of new computation of Hel emissivity Porter, Bauman, Ferland, MacAdam 2006

PBFM

- Peimbert, Luridiana & Peimbert 2007
  - 5 HII regions of IT98
  - use Hel emission lines to determine T

 $Y_p = 0.249 \pm 0.009$ 

- Izotov, Thuan & Stasinska 2007
  - 93 HII regions (HeBCD) + 271 HII regions in SDSS DR5
  - T(HeII) = (0.95 1.0)T(OIII)
  - underlying stellar absorption

 $Y_p = 0.2516 \pm 0.0011$ 

# New Emissivity



# Izotov, Thuan, Stasinska 2007





# **PBFM** $Y_p = 0.2516 \pm 0.0011$

# Systematic errors

- He I emissivity
- T(OIII) may be different from T(HeII)
- Underlying Hel stellar absorption
- Collisional excitation of hydrogen emission lines
- Hell and Hll regions may not coincident

correction factor  $ICF(He^+ + He^{2+})$ 

# Error Budget IT (2007)

Property	$\Delta Y_p$
He I emissivity	$\lesssim +1.7\%$
$T_e({\rm He^+}) = (0.95 - 1.0) \times T_e({\rm O~{III}})$	$\lesssim -1.0\%$
Underlying He I stellar absorption	$\lesssim +3.0\%$
Collisional excitation of hydrogen emission lines	$\lesssim +1.0\%$
$ICF(\mathrm{He^{+} + He^{2+}})$	$\lesssim -1.0\%$

# Yp History



# 3. Li7



Spite plateau [Spite & Spite (1987)]

constant Li7 abundance in warmest metal-poor stars



Primordial abundance of Li 7



Bonifacio, Molaro 1997

# 6708Å line



# Recent works

- Bonifacio & Malaro (1997)
  - 41 metal-poor stars
  - IRFM to determine T
  - no dep. on [Fe/H]

T is found by comparison of infrared flux with bolometric flux

 $\log_{10}({^7\text{Li}/\text{H}})_p = -9.762 \pm 0.012(sta) \pm 0.05(sys)$ 

- Ryan et al. (2000)
  - 23 metal-poor field stars
  - IRFM

• correlation between Li and [Fe/H]  $\log_{10}(^{7}\text{Li}/\text{H})_{p} = -9.91 \pm 0.10$ 

## Infrared flux method

• Effective Temperature

$$\sigma T_{eff}^4 = F = f(r/R)^2$$

F: surface fluxf: observed fluxr: distance to starR: Radius of star

• Monochromatic flux

 $(r/R)^2 = F(\lambda)/f(\lambda)$ 

 $f(\lambda)$ : observed monochromatic flux  $F(\lambda)$ : monochromatic surface flux

$$\sigma T_{eff}^4 = fF(\lambda)/f(\lambda)$$

model atmosphere calc.

 Infrared wavelength is used because the Planck curve is only weakly dependent at infrared wavelength, and hence small uncertainty in choice of a model atmosphere

## $\log_{10}(^{7}\text{Li/H}) = (-9.95 \pm 0.05) + (0.118 \pm 0.023)$ [Fe/H]



Ryan et al (2000)

# Recent works (cont.)

- Bonifacio et al. (2002)
  - 12 stars in metal-poor globular cluster NGC6397

$$\log_{10}(^{7}\text{Li/H})_{p} = -9.66 \pm 0.056$$

[Fe/H] = -2.03

- Melendez & Ramirez (2004)
  - 41 metal-poor dwarf stars
  - new calibration of IRFM \_\_\_\_\_



higher Li abundance

no correlation between Li and [Fe/H] 

 $\log_{10}({^7\text{Li}/\text{H}})_p = -9.63 \pm 0.06$ 



Melendez & Ramirez (2000)

# Recent works (cont.)

- Asplund et al. (2005)
  - 24 metal-poor halo dwarfs
  - $H\alpha$  line profile to determine T
  - correlation between Li and [Fe/H]

 $\log_{10}({^7\text{Li}/\text{H}})_p = -9.90 \pm 0.06$ 

- Bonifacio et al (2006)
  - I9 metal-poor dwarf stars
  - $H\alpha$  line profile to determine T
  - no correlation between Li and [Fe/H]





**rig. A.1.** Example of fits to the H $\alpha$  line of the star BS 16023-043. The best-fit profile corresponds to  $T_{\text{eff}} = 6364$  K. The ot we profiles shown correspond to  $T_{\text{eff}} \pm 200$  K of this value. The narrow absorption features are H<sub>2</sub>O telluric lines.



 $\log_{10}({^7\text{Li}/\text{H}}) = (-9.59 \pm 0.02) + (0.103 \pm 0.010)[\text{Fe}/\text{H}]$ 

Asplund et al (2005)



Bonifacio et al (2006)

# Effect of different temperature scales



Bonifacio et al (2006)

Melendez & Ramirez (2000)

# Lithium Problem

- WMAP Prediction  $\log_{10}(^{7}\text{Li}/\text{H})_{p} = -9.35 \pm 0.10$
- Observation large discrepancy log<sub>10</sub>(<sup>7</sup>Li/H)<sub>p</sub> = −9.90 ± 0.10
   Dpletion ? rotational mixing → at most D<sub>7</sub> = 0.3 dex

Pinsonneault et al (2002)

## Rotational mixing and Li7 abundance

Pinsonneault et al (2002)

- Rotation induces mixing in the radiative interiors of stars, leading to surface Li depletion during mainsequence phase
- Ryan, Norris & Beers (1999) sample is fully consistent with mild rotational mixing induced depletion

$$D_7 = 0.2 \pm 0.1$$





# Li6

- Asplund et al (2005)
  - Li6 was detected in 9 out of 24 metal-poor halo dwarfs
  - Detection of Li6 in very metal-poor star LP 815-43

 $^{6}\mathrm{Li}/^{7}\mathrm{Li} = 0.046 \pm 0.022$ 

 $[\mathrm{Fe}/\mathrm{H}] = -2.74$ 

This Li6 abundance may be primordial

# Detection of Li6



Asplund et al (2005)

# Detection of Li6



### Asplund et al (2005)

# Li6

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    SBBN prediction
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 $^{6}\mathrm{Li}/^{7}\mathrm{Li} \simeq 3 \times 10^{-5}$ 

• Depletion ?

 $D_6 \simeq 2.5 D_7$  $\log_{10}(^6 \text{Li}/^7 \text{Li})_p = 1.5 D_7 + \log_{10}(^6 \text{Li}/^7 \text{Li})_{obs}$ 

# Chemical evolution model

- spallation process (p + O, C, N)
- $\alpha$ - $\alpha$  fusion reactions



5. D

D

- Absorption lines in Damped Lyα systems along sight lines of QSOs
  - Burles & Tytler (1998) PKS 1937-1009 (z=3.572) Q 1009+299 (z=2.504) O'Meara et al (2001) HS 0105+1619 (z=2.536) Kirkman et al (2003) <u>Q |243+3047</u> (z=2.252) • O'Meara et al (2006) SDSSI558-0031 (z=2.702)
  - Pettini & Bowen (2001)
     Q 2206-199 (z=2.076)

 $D/H = (3.25 \pm 0.3) \times 10^{-5}$  $D/H = 3.98^{+0.59}_{-0.67} \times 10^{-5}$ 

 $D/H = (2.54 \pm 0.23) \times 10^{-5}$ 

$$D/H = 2.42^{+0.35}_{-0.25} \times 10^{-5}$$

 $D/H = 2.88^{+0.49}_{-0.43} \times 10^{-5}$ 

 $D/H = (1.65 \pm 0.35) \times 10^{-5}$ 

## D absorption in QSO spectrum





 $F_{\lambda} \times 10^{-16} \text{ (ergs sec}^{-1} \text{ cm}^{-2} \text{ Å}^{-1}\text{)}$ 

# D/H vs N<sub>HI</sub>



No plausible mechanism to explain correlation

# D/H vs N<sub>HI</sub>



However, a single value for D/H is still not supported

# D/H vs N<sub>HI</sub>

#### weighted mean



However, a single value for D/H is still not supported

# Conclusion

- 元素合成の理論と観測はよく合っているいる
   が、より精密な定量的比較を行うにはもっと
   系統誤差の理解が必要
- 宇宙のバリオン密度を精度良く決める役割は CMBにとられた
- しかし、現在でも宇宙の最も初期を探ることのできる重要なプローブである
    $N_{\nu}$

constraints on exotic particles



# Extinction Law Comparison

Whitford (1958) as parameterized by Miller & Mathews 1972 and Izotov, Thuan & Lipovetsky 1994



# **Extinction Curve**





Decrease with increasing temp is due to decreasing population of 2<sup>3</sup>S by collisional excitation

## fitting : $f(\tau) = f_0 + f_1 \tau + f_2 \tau^2$

Benjamin, Skillman, Smits 2002, ApJ 569,288



# Fitting formula $f(\tau) =$ $1 + (\tau/2)[a + (b_0 + b_1 n_e + b_2 n_2^2)T]$ $(\tau \le 2.0)$

## large difference from IZ 98

Benjamin, Skillman, Smits 2002, ApJ 569,288



# New Emissivity



# Optical depth functions

$$f(3889) = 1 + (\tau/2) \left[ -0.106 + (5.14 \times 10^{-5} - 4.20 \times 10^{-7} n_e + 1.97 \times 10^{-10} n_e^2) T \right]$$

 $f(4026) = 1 + (\tau/2) \left[ 0.00143 + (4.05 \times 10^{-4} + 3.63 \times 10^{-8} n_e)T) \right]$ 

- $f(4471) = 1 + (\tau/2) \left[ 0.00274 + (8.81 \times 10^{-4} 1.21 \times 10^{-6} n_e)T) \right]$
- $f(5876) = 1 + (\tau/2) \left[ 0.00470 + (2.23 \times 10^{-3} 2.51 \times 10^{-6} n_e)T) \right]$
- f(6678) = 1
- $f(7065) = 1 + (\tau/2) \left[ 0.359 + (-3.46 \times 10^{-2} 1.84 \times 10^{-4} n_e + 3.039 \times 10^{-7} n_e^2) T \right]$

$$au = au_{3889} = n(2^3S)\sigma_{3889}R_S$$
  
 $n(2^3S)$  :density of HeI in the metastable state  
 $R_S$  :Stromgren radius

#### Crighton, Webb, Ortiz-Gil, Fernandez-Soto (2004)

